## Geometry of extended Bianchi-Cartan-Vranceanu spaces

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#### Abstract

The differential geometry of 3-dimensional Bianchi, Cartan and Vranceanu (BCV) spaces is well known. We introduce the extended Bianchi, Cartan and Vranceanu (EBCV) spaces as a natural seven dimensional generalization of BCV spaces and study some of their main geometric properties, such as the Levi-Civita connection, Ricci curvatures, Killing fields and geodesics.

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# 1 The Bianchi-Cartan-Vranceanu (BCV) spaces (see [2, 5])

It was Cartan ([6]) who obtained the families of today known as BCV-spaces by classifying three-dimensional Riemannian manifolds with four-dimensional isometry group. They also appeared in the work of L. Bianchi ([3, 4]), and G. Vranceanu ([18]). These kind of spaces have been extensively studied and classified (see for instance [14, 17]). In theoretical cosmology they are known as Bianchi-Kantowski-Saks spaces, which are used to construct some homogeneous spacetimes ([10]).

For real numbers m and l, consider the set

$$BCV(m,l) = \{(x,y,z) \in \mathbb{R}^3 : 1 + m(x^2 + y^2) > 0\}$$

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equipped with the metric

$$ds_{m,l}^2 = \frac{dx^2 + dy^2}{\lambda^2} + \left(dz - \frac{l}{2} \frac{xdy - ydx}{\lambda}\right)^2,$$

where m, l are real numbers and  $\lambda = 1 + m(x^2 + y^2) > 0$ .

Observe that this metric is obtained as a conformal deformation of the planar Euclidean metric by adding the imaginary part of  $z d\bar{z}$ , for a complex number z.

The complete classification of BCV spaces is as follows:

- (i) If m = l = 0, then  $BCV(m, l) \cong \mathbb{R}^3$ ;
- (ii) If  $m = \frac{l}{4}$ , then  $BCV(m, l) \cong (\mathbb{S}^3(m) \{\infty\})$ ;
- (iii) If m > 0 and l = 0, then  $BCV(m, l) \cong (\mathbb{S}^2(4m) \{\infty\}) \times \mathbb{R}$ ;
- (iv) If m < 0 and l = 0, then  $BCV(m, l) \cong (\mathbb{H}^2(4m) \{\infty\}) \times \mathbb{R}$ ;
- (v) If m > 0 and  $l \neq 0$ , then  $BCV(m, l) \cong SU(2) {\infty}$ ;
- (vi) If m < 0 and  $l \neq 0$ , then  $BCV(m, l) \cong \widetilde{SL}(2, \mathbb{R})$ ;
- (vii) If m = 0 and  $l \neq 0$ , then  $BCV(m, l) \cong Nil_3$ .

The following vector fields form an orthonormal frame of BCV(m, l):

$$E_1 = \lambda \partial_x - \frac{l}{2}y\partial_z, \qquad E_2 = \lambda \partial_y + \frac{l}{2}x\partial_z, \qquad E_3 = \partial_z.$$

Let  $\mathcal{D}$  be the distribution generated by  $\{E_1, E_2\}$ , then the manifold  $(BCV(m, l), \mathcal{D}, ds_{m,l}^2)$  is an example of sub-Riemannian geometry (see [5, 15]) and the horizontal distribution is a 2-step breaking-generating distribution everywhere.

## 2 Extended Bianchi-Cartan-Vranceanu spaces

#### 2.1 Set up

Observe that letting z = x + iy, we see that  $\text{Im}(z d\bar{z}) = ydx - xdy$ , which reminds us the map  $\mathbb{C} \times \mathbb{C} \to \mathbb{R} \times \mathbb{C}$  given by  $(z_1, z_2) \mapsto (|z_1|^2 - |z_2|^2, 2(z_1\bar{z}_2))$ , that easily leads to the classical Hopf fibration  $\mathbb{S}^1 \hookrightarrow \mathbb{S}^3 \to \mathbb{S}^2$ , where coordinates in  $\mathbb{S}^2$  are given by  $(|z_1|^2 - |z_2|^2, 2\text{Re}(z_1\bar{z}_2), 2\text{Im}(z_1\bar{z}_2))$ .

In the same line, using quaternions  $\mathbb{H}$  instead of complex numbers, we get the fibration  $\mathbb{S}^3 \hookrightarrow \mathbb{S}^7 \to \mathbb{S}^4$ . Quaternions are usually presented with the imaginary units i, j, k in the form  $q = x_0 + x_1 i + x_2 j + x_3 k$ ,  $x_0, x_1, x_2, x_3 \in \mathbb{R}$  with  $i^2 = j^2 = k^2 = ijk = -1$ . They can also be defined equivalently, using the complex numbers  $c_1 = x_0 + x_1 i$  and  $c_2 = x_2 + x_3 i$ , in the form  $q = c_1 + c_2 j$ . Then for a point  $(q_1 = \alpha + \beta j, q_2 = \gamma + \delta j) \in \mathbb{S}^7$ , we get the following coordinate expressions  $(|q_1|^2 - |q_2|^2, 2\text{Re}(\bar{\alpha}\gamma + \bar{\beta}\delta), 2\text{Im}(\bar{\alpha}\gamma + \bar{\beta}\delta), 2\text{Re}(\alpha\delta - \beta\gamma), 2\text{Im}(\alpha\delta - \beta\gamma))$ .

For any  $q = w + xi + yj + zk \in \mathbb{H}$  we find that  $qd\bar{q} = wdw + xdx + ydy + zdz + (xdw - wdx + zdy - ydz)i + (ydw - wdy + xdz - zdx)j + (zdw - wdz + ydx - xdy)k$ . The quaternionic contact

group  $\mathbb{H} \times \text{Im}\mathbb{H}$ , with coordinates (w, x, y, z, r, s, t), can be equipped with the metric

$$ds^{2} = (dw^{2} + dx^{2} + dy^{2} + dz^{2}) + \left(dr + \frac{1}{2}(xdw - wdx + zdy - ydz)\right)^{2} + \left(ds + \frac{1}{2}(ydw - wdy + xdz - zdx)\right)^{2} + \left(dt + \frac{1}{2}(zdw - wdz + ydx - xdy)\right)^{2}.$$

Then, by extending this metric, and following [9], it seems natural to find a 7-dimensional generalization of the 3-dimensional BCV spaces endowed with the two-parameter family of metrics

$$ds_{m,l}^{2} = \frac{dw^{2} + dx^{2} + dy^{2} + dz^{2}}{K^{2}} + \left(dr + \frac{l}{2}\frac{wdx - xdw + ydz - zdy}{K}\right)^{2} + \left(ds + \frac{l}{2}\frac{wdy - ydw + zdx - xdz}{K}\right)^{2} + \left(dt + \frac{l}{2}\frac{wdz - zdw + xdy - ydx}{K}\right)^{2},$$

where m, l are real numbers and  $K = 1 + m(w^2 + x^2 + y^2 + z^2) > 0$ .

Then  $(EBCV, ds_{m,l}^2)$  will be called extended BCV spaces (EBCV for short).

Note that the first summand in the metric  $ds_{m,l}^2$  is as a conformal change of the Euclidean metric on  $\mathbb{R}^4$ , whereas the three other summands (depending on m and l) are coming from the imaginary part of  $q d\bar{q}$ , for a quaternion q. When m=0 we get a one-parameter family of Riemannian metrics depending on l. Furthermore, if l=1, we find the 7-dimensional quaternionic Heisenberg group (see [9] and [19]). The manifold EBCV provides another example of sub-Riemannian geometry and the horizontal distribution is a 2-step breaking-generating distribution everywhere.

Observe that when m=l=0, EBCV is nothing but  $\mathbb{R}^7$ ; when m>0, l=0,  $EBCV\cong \mathbb{S}^4(4m)\times\mathbb{R}^3$  and when m<0, l=0,  $EBCV\cong \mathbb{H}^4(4m)\times\mathbb{R}^3$ .

The metric  $ds_{m,l}^2$  can also be written as

$$\mathrm{d} \mathrm{s}_{m,l}^2 = \sum_{\alpha=1}^7 \omega^\alpha \otimes \omega^\alpha,$$

where

$$\begin{split} \omega^1 &= dr + \frac{l}{2K}(wdx - xdw + ydz - zdy), & \omega^4 &= \frac{1}{K}dw, \\ \omega^2 &= ds + \frac{l}{2K}(wdy - ydw + zdx - xdz), & \omega^5 &= \frac{1}{K}dx, \\ \omega^3 &= dt + \frac{l}{2K}(wdz - zdw + xdy - ydx), & \omega^6 &= \frac{1}{K}dy, \\ \omega^7 &= \frac{1}{K}dz, \end{split}$$

with the corresponding dual orthonormal frame

$$X_1 = \partial_r, \qquad X_2 = \partial_s, \qquad X_3 = \partial_t,$$

$$X_{4} = K\partial_{w} + \frac{lx}{2}\partial_{r} + \frac{ly}{2}\partial_{s} + \frac{lz}{2}\partial_{t}, \qquad X_{5} = K\partial_{x} - \frac{lw}{2}\partial_{r} - \frac{lz}{2}\partial_{s} + \frac{ly}{2}\partial_{t},$$

$$X_{6} = K\partial_{y} + \frac{lz}{2}\partial_{r} - \frac{lw}{2}\partial_{s} - \frac{lx}{2}\partial_{t}, \qquad X_{7} = K\partial_{z} - \frac{ly}{2}\partial_{r} + \frac{lx}{2}\partial_{s} - \frac{lw}{2}\partial_{t}.$$

Writing  $1 \le i, j \le 3, 4 \le a \le 7$ , we find that

$$[X_i, X_j] = 0;$$
  $[X_i, X_a] = 0,$ 

as well as

 $[X_4, X_5] = -l\{1 + m(y^2 + z^2)\}X_1 + ml(wz + xy)X_2 - ml(wy - xz)X_3 - 2mxX_4 + 2mwX_5,$ and so on (see Appendix).

For later use, when m = 0 brackets reduce to

$$[X_4, X_5] = -lX_1,$$
  $[X_4, X_6] = -lX_2,$   $[X_4, X_7] = -lX_3,$   $[X_5, X_6] = -lX_3,$   $[X_5, X_7] = lX_2,$   $[X_6, X_7] = -lX_1.$ 

**Remark 1** When l = 1, we have the brackets of the quaternionic contact manifold.

As for the Levi-Civita connection we find out

$$\nabla_{X_i} X_j = 0, \qquad \nabla_{X_i} X_a = \nabla_{X_a} X_i,$$

and

$$\nabla_{X_1} X_4 = \frac{l}{2} \{1 + m(y^2 + z^2) X_5 + \frac{ml}{2} (wz - xy) X_6 - \frac{ml}{2} (wy + xz) X_7,$$

$$\nabla_{X_1} X_5 = -\frac{l}{2} \{1 + m(y^2 + z^2) \} X_4 + \frac{ml}{2} (wy + xz) X_6 + \frac{ml}{2} (wz - xy) X_7,$$

$$\nabla_{X_1} X_6 = -\frac{ml}{2} (wz - xy) X_4 - \frac{ml}{2} (wy + xz) X_5 + \frac{l}{2} \{1 + m(w^2 + x^2) \} X_7,$$

$$\nabla_{X_1} X_7 = \frac{ml}{2} (wy + xz) X_4 - \frac{ml}{2} (wz - xy) X_5 - \frac{l}{2} \{1 + m(w^2 + x^2) \} X_6,$$

and son on (see Appendix).

When m=0, the Levi-Civita connection reduces to

$$\begin{array}{llll} \nabla_{X_1} X_4 = \frac{l}{2} X_5, & \nabla_{X_3} X_4 = \frac{l}{2} X_7, & \nabla_{X_5} X_4 = \frac{l}{2} X_1, & \nabla_{X_7} X_4 = \frac{l}{2} X_3, \\ \nabla_{X_1} X_5 = -\frac{l}{2} X_4, & \nabla_{X_3} X_5 = \frac{l}{2} X_6, & \nabla_{X_5} X_5 = 0, & \nabla_{X_7} X_5 = -\frac{l}{2} X_2, \\ \nabla_{X_1} X_6 = \frac{l}{2} X_7, & \nabla_{X_3} X_6 - \frac{l}{2} X_5, & \nabla_{X_5} X_6 = -\frac{l}{2} X_3, & \nabla_{X_7} X_6 = \frac{l}{2} X_1, \\ \nabla_{X_1} X_7 = -\frac{l}{2} X_6, & \nabla_{X_3} X_7 = -\frac{l}{2} X_4, & \nabla_{X_5} X_7 = \frac{l}{2} X_2, & \nabla_{X_7} X_7 = 0. \\ \nabla_{X_2} X_4 = \frac{l}{2} X_6, & \nabla_{X_4} X_4 = 0, & \nabla_{X_6} X_4 = \frac{l}{2} X_2, & \nabla_{X_7} X_7 = 0. \\ \nabla_{X_2} X_5 = -\frac{l}{2} X_7, & \nabla_{X_4} X_5 = -\frac{l}{2} X_1, & \nabla_{X_6} X_5 = \frac{l}{2} X_3, \\ \nabla_{X_2} X_6 = -\frac{l}{2} X_4, & \nabla_{X_4} X_6 = -\frac{l}{2} X_2, & \nabla_{X_6} X_6 = 0, \\ \nabla_{X_2} X_7 = \frac{l}{2} X_5, & \nabla_{X_4} X_7 = -\frac{l}{2} X_3, & \nabla_{X_6} X_7 = -\frac{l}{2} X_1, \end{array}$$

**Remark 2** When l = 1, we find the Levi-Civita connection of the quaternionic contact manifold.

As for the curvature tensor R we have

$$R_{X_1X_4X_1X_4} = R_{X_1X_5X_1X_5} = \frac{l^2}{4} \{ 1 + m(K+1)(y^2 + z^2) \},$$
  

$$R_{X_1X_6X_1X_6} = R_{X_1X_7X_1X_7} = \frac{l^2}{4} \{ 1 + m(K+1)(w^2 + x^2) \},$$

and so on (see Appendix).

**Remark 3** When m = 0, the curvature of the quaternionic contact manifold reduces to

$$R_{X_i X_a X_i X_a} = \frac{l^2}{4}; \quad R_{X_a X_b X_a X_b} = -\frac{3l^2}{4}.$$

#### 2.2 The Ricci tensor

**Proposition 4** The matrix representing the Ricci tensor is given by

$$\begin{pmatrix} \frac{l^2}{2}(K^2+1) & 0 & 0 & \vdots \\ 0 & \frac{l^2}{2}(K^2+1) & 0 & \vdots \\ 0 & 0 & \frac{l^2}{2}(K^2+1) & \vdots \\ -mlx(K+2) & -mly(K+2) & -mlz(K+2) & \vdots \\ mlw(K+2) & mlz(K+2) & -mly(K+2) & \vdots \\ -mlz(K+2) & mlw(K+2) & mlx(K+2) & \vdots \\ mly(K+2) & -mlx(K+2) & mlw(K+2) & \vdots \\ \vdots & -mlx(K+2) & mlw(K+2) & \vdots \\ \vdots & -mlx(K+2) & mlx(K+2) & mlw(K+2) & -mlx(K+2) \\ \vdots & -mly(K+2) & mlz(K+2) & mlw(K+2) & -mlx(K+2) \\ \vdots & -mlz(K+2) & -mly(K+2) & mlx(K+2) & mlw(K+2) \\ \vdots & -mlz(K+2) & -mly(K+2) & mlx(K+2) & mlw(K+2) \\ \vdots & -mlz(K+2) & -mly(K+2) & mlx(K+2) & mlw(K+2) \\ \vdots & -mlz(K+2) & -mly(K+2) & mlx(K+2) & mlw(K+2) \\ \vdots & -mlz(K+2) & -mly(K+2) & mlx(K+2) & mlw(K+2) \\ \vdots & -mlz(K+2) & -mly(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & -mly(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & -mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) \\ \vdots & -mlx(K+2) & mlx(K+2) & mlx(K+2) \\ \vdots & -ml$$

where  $A = -l^2(K+1)$  and  $B = 12m - 3/2l^2$ .

Some particular cases could be interesting, for instance we get the following Ricci matrix when K = 1 (or m = 0)

$$\operatorname{Ric}_{1} = \begin{pmatrix} l^{2} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & l^{2} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & l^{2} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -3/2l^{2} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -3/2l^{2} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -3/2l^{2} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3/2l^{2} \end{pmatrix}$$

**Remark 5** When l = 1, we find the Ricci curvature of the quaternionic contact manifold.

An easy computation leads to

Corollary 6 The EBCV manifold has constant scalar curvature S = 48m.

## 3 The characteristic connection on the EBCV manifold

We consider on EBCV the characteristic connection D defined by (see [7])

$$D_L M = \nabla_L M + \frac{P}{2} (\nabla_L P) M,$$

where P is the natural almost product structure given by  $P = \mathcal{V} - \mathcal{H}$ ,  $Id = \mathcal{V} + \mathcal{H}$  and L, M are arbitrary vector fields. Let us remember that the vertical distribution in EBCV is spanned by  $X_1, X_2, X_3$  and the horizontal distribution by  $X_4, X_5, X_6, X_7$ . Then we have

$$D_{X_{i}}X_{j} = \mathcal{V}(\nabla_{X_{i}}X_{j}), i, j = 1, 2, 3,$$

$$D_{X_{a}}X_{j} = \mathcal{V}(\nabla_{X_{a}}X_{j}), a = 4 \dots, 7; j = 1, 2, 3,$$

$$D_{X_{i}}X_{b} = \mathcal{H}(\nabla_{X_{i}}X_{b}), i = 1, 2, 3; b = 4, \dots, 7,$$

$$D_{X_{a}}X_{b} = \mathcal{H}(\nabla_{X_{a}}X_{b}), a, b = 4, \dots, 7.$$

This is a metric connection which can be completely obtained by using the table giving the Levi-Civita connection.

Following the classification given by A. M. Naveira for almost product structures, [12], we have

**Proposition 7** (EBCV, P) is in (TGF, AF) class.

To prove this proposition it is enough to see that  $\nabla_A(P)B = 0$ , when A, B are vertical, and  $\nabla_X(P)X = 0$ , if X is horizontal. The result follows using the tables given in the Appendix for the Levi-Civita connection.

When we consider m=0, it is known that EBCV is a homogeneous manifold. Indeed, it is the quaternionic contact group (see [9, 19]). In [1] W. Ambrose and I. M. Singer proved that a connected, complete and simply-connected Riemannian manifold (M,g) is homogeneous if and only if there exists a (1,2) tensor field T such that

(i) 
$$g(T_XY, Z) + g(Y, T_XZ) = 0$$
,

(ii) 
$$(\nabla_X R)_{YZ} = [T_x, R_{YZ}] - R_{T_X YZ} - R_{YT_X Z},$$

(iii) 
$$(\nabla_X T)_Y = [T_X, T_Y] - T_{T_X Y},$$

for  $X, Y, Z \in \mathfrak{X}(M)$ , where  $\nabla$  stands for the Levi-Civita connection and R is the Riemann curvature tensor of M (see [16]). As a consequence, Tricerri and Vanhecke define a homogeneous Riemannian structure on (M, g) as a (1,2) tensor field T which is a solution of the above three equations. Instead of taking (1,2) tensors it is preferred to work with (0,3) tensors via the isomorphism  $T_{uvw} = g(T_uv, w)$ , for  $u, v, w \in T_pM$  and  $p \in M$ . So far we have not been able to find the tensor field T satisfying the above conditions, which are equivalent to those given in page 14 of [16]. Then allow us to state the following question: is EBCV a homogeneous manifold?

Let  $T = T^D$  be the torsion tensor of the connection D, that is,

$$T_L^D M \equiv T^D(L, M) = D_L M - D_M L - [L, M],$$

or equivalently

$$T^{D}(L, M) = \frac{P}{2} \left( (\nabla_{L} P) M - (\nabla_{M} P) L \right).$$

Then we find out

$$T^D(X_k, X_k) = 0, \ k = 1, \dots 7.$$

There are non-vanishing components such as  $T^D(X_i, X_a)$  or  $T^D(X_a, X_b)$ , for instance,

$$T^{D}(X_{1}, X_{4}) = \frac{l}{2} \{1 + m(y^{2} + z^{2})\}X_{5} + \frac{lm}{2}(wz - xy)X_{6} - \frac{lm}{2}(wy + xz)X_{7}$$

or

$$T^{D}(X_4, X_5) = l\{(1 + m(y^2 + z^2))X_1 - m(wz + xy)X_2 + m(wy - xz)X_3\}.$$

On the other hand, it is easy to see that

(a) 
$$T_{X_1X_4X_5}^D + T_{X_5X_1X_4}^D + T_{X_4X_5X_1}^D = \langle T_{X_1}^D X_4, X_5 \rangle + \langle T_{X_5}^D X_1, X_4 \rangle + \langle T_{X_4}^D X_5, X_1 \rangle = 2l \neq 0;$$

(b) 
$$T_{XYZ}^D + T_{YXZ}^D = 0.$$

### 4 Killing vector fields in EBCV

Remember that a Killing vector field is a vector field on a Riemannian manifold that preserves the metric. Killing vector fields are the infinitesimal generators of isometries, that is, flows generated by Killing fields are continuous isometries of the manifold. Specifically, a vector field X is a Killing vector field if the Lie derivative with respect to X of the metric X or equivalently

$$\mathcal{L}_X \mathrm{ds}_{l,m}^2 = (\mathcal{L}_X \omega^\alpha) \otimes \omega^\alpha = 0, \tag{1}$$

where

$$\mathcal{L}_X \omega^\alpha = \iota_X d\omega^\alpha + d(\iota_X \omega^\alpha).$$

In terms of the Levi-Civita connection, Killing's condition is equivalent to

$$g(\nabla_Y X, Z) + g(Y, \nabla_Z X) = 0. (2)$$

It is easy to prove that

**Proposition 8**  $\mathcal{L}_X g(Y, Z) = 0$  if and only if  $\mathcal{L}_X g(X_i, X_j) = 0$  for basic vector fields  $X_i, X_j$ .

We know that the dimension of the Lie algebra of the Killing vector fields is  $m \le n(n+1)/2$  and the maximum is reached on constant curvature manifolds ([8], p. 238, Vol. II), then for our manifold m < 28. Then obviously

**Proposition 9** The basic vertical vector fields  $X_1, X_2, X_3$  are Killing fields.

From (2) it is easy to prove that the horizontal basic vector fields  $X_4, \ldots, X_7$  are not Killing vector fields.

In [13] the Levi-Civita connection, curvature tensor and Killing vector fields on Bianchi-Cartan-Vranceanu spaces are introduced. In [14] Piu and Profir proved that the Lie algebra of Killing vector fields of BCV spaces is 4-dimensional for generic parameters m and l.

Now we are going to determine the space of Killing vector fields in EBCV.

#### 4.1 The Killing equations

In the usual coordinate system (r, s, t, w, x, y, z) on EBCV, a vector field  $X = \sum_{\alpha=1}^{7} f_{\alpha} X_{\alpha}$  will be a Killing field if and only if the real functions  $f_i$  satisfy the following system of 28-partial differential equations:

$$\begin{split} &\partial_r(f_1) = 0, \\ &\partial_s(f_2) = 0, \\ &\partial_t(f_3) = 0, \\ &\partial_r(f_2) + \partial_s(f_1) = 0, \\ &\partial_r(f_3) + \partial_t(f_1) = 0, \\ &\partial_r(f_3) + \partial_t(f_2) = 0, \\ &\partial_r(f_3) + \mathcal{K}\partial_w(f_1) + \frac{t_2}{2}\partial_s(f_1) + \frac{t_2}{2}\partial_t(f_1) - l\{1 + m(y^2 + z^2)\}f_3 - ml(wz - xy)f_6 + ml(wy + xz)f_7 = 0, \\ &\partial_r(f_3) + \mathcal{K}\partial_w(f_1) + \frac{t_2}{2}\partial_s(f_1) + \frac{t_2}{2}\partial_t(f_1) + l\{1 + m(y^2 + z^2)\}f_3 - ml(wz - xy)f_6 + ml(wz - xy)f_7 = 0, \\ &\partial_r(f_3) + \mathcal{K}\partial_w(f_1) - \frac{t_2}{2}\partial_s(f_1) + \frac{t_2}{2}\partial_t(f_1) + ml(wz - xy)f_4 + ml(wy + xz)f_5 - ml(wz - xy)f_7 = 0, \\ &\partial_r(f_6) + \mathcal{K}\partial_y(f_1) - \frac{t_2}{2}\partial_s(f_1) - \frac{t_2}{2}\partial_t(f_1) + ml(wz - xy)f_4 + ml(wy + xz)f_5 - l\{1 + m(w^2 + x^2)f_7 = 0, \\ &\partial_r(f_7) + \mathcal{K}\partial_s(f_1) + \frac{t_2}{2}\partial_s(f_1) - \frac{t_2}{2}\partial_t(f_1) - ml(wy + xz)f_4 + ml(wz - xy)f_5 + l\{1 + m(w^2 + x^2)f_6 = 0, \\ &\partial_s(f_4) + \mathcal{K}\partial_w(f_2) + \frac{t_2}{2}\partial_r(f_2) + \frac{t_2}{2}\partial_t(f_2) + ml(wz + xy)f_5 - l\{1 + m(x^2 + z^2)\}f_6 - ml(wx - yz)f_7 = 0, \\ &\partial_s(f_3) + \mathcal{K}\partial_x(f_2) - \frac{t_2}{2}\partial_r(f_2) + \frac{t_2}{2}\partial_t(f_2) - ml(wz + xy)f_4 + ml(wx - yz)f_6 + l\{1 + m(w^2 + y^2)\}f_7 = 0, \\ &\partial_s(f_6) + \mathcal{K}\partial_y(f_2) - \frac{t_2}{2}\partial_r(f_2) - \frac{t_2}{2}\partial_r(f_2) + l\{1 + m(x^2 + z^2)\}f_4 - ml(wx - yz)f_5 - ml(wz + xy)f_7 = 0, \\ &\partial_s(f_7) + \mathcal{K}\partial_s(f_2) - \frac{t_2}{2}\partial_r(f_2) - \frac{t_2}{2}\partial_r(f_2) + ml(wx - yz)f_4 - l\{1 + m(w^2 + y^2)\}f_5 + ml(wz + xy)f_6 = 0, \\ &\partial_t(f_4) + \mathcal{K}\partial_w(f_3) + \frac{t_2}{2}\partial_r(f_3) + \frac{t_2}{2}\partial_s(f_3) + ml(wy - xz)f_5 + ml(wx + yz)f_6 - l\{1 + m(x^2 + y^2)\}f_7 = 0, \\ &\partial_t(f_3) + \mathcal{K}\partial_s(f_3) + \frac{t_2}{2}\partial_r(f_3) - \frac{t_2}{2}\partial_s(f_3) + ml(wy - xz)f_3 + l\{1 + m(w^2 + z^2)\}f_6 + ml(wx + yz)f_7 = 0, \\ &\partial_t(f_6) + \mathcal{K}\partial_y(f_3) + \frac{t_2}{2}\partial_r(f_3) - \frac{t_2}{2}\partial_s(f_3) + ml(wy - xz)f_3 + l\{1 + m(w^2 + z^2)\}f_6 + ml(wx + yz)f_7 = 0, \\ &\partial_t(f_6) + \mathcal{K}\partial_y(f_3) + \frac{t_2}{2}\partial_r(f_3) + \frac{t_2}{2}\partial_r(f_3)$$

It seems that the solution of the system is very difficult, so that we focus on solving the system for m = 0, that is:

$$\begin{split} &\partial_r(f_1) = 0, \\ &\partial_s(f_2) = 0, \\ &\partial_t(f_3) = 0, \\ &\partial_r(f_2) + \partial_s(f_1) = 0, \\ &\partial_r(f_3) + \partial_t(f_1) = 0, \\ &\partial_r(f_3) + \partial_t(f_1) = \frac{1}{2}\partial_s(f_1) + \frac{1}{2}\partial_t(f_1) - lf_5 = 0, \\ &\partial_r(f_3) + \partial_t(f_1) - \frac{1}{2}\partial_s(f_1) + \frac{1}{2}\partial_t(f_1) + lf_4 = 0, \\ &\partial_r(f_3) + \partial_t(f_1) - \frac{1}{2}\partial_s(f_1) - \frac{1}{2}\partial_t(f_1) + lf_6 = 0, \\ &\partial_r(f_3) + \partial_x(f_1) - \frac{1}{2}\partial_s(f_1) - \frac{1}{2}\partial_t(f_1) + lf_6 = 0, \\ &\partial_r(f_7) + \partial_z(f_1) + \frac{1}{2}\partial_s(f_1) - \frac{1}{2}\partial_t(f_2) + lf_6 = 0, \\ &\partial_r(f_7) + \partial_z(f_1) + \frac{1}{2}\partial_s(f_1) - \frac{1}{2}\partial_t(f_2) + lf_7 = 0, \\ &\partial_s(f_3) + \partial_w(f_2) + \frac{1}{2}\partial_r(f_2) + \frac{1}{2}\partial_t(f_2) + lf_7 = 0, \\ &\partial_s(f_3) + \partial_w(f_2) + \frac{1}{2}\partial_r(f_2) + \frac{1}{2}\partial_t(f_2) + lf_4 = 0, \\ &\partial_s(f_3) + \partial_x(f_2) - \frac{1}{2}\partial_r(f_2) - \frac{1}{2}\partial_t(f_2) + lf_4 = 0, \\ &\partial_s(f_3) + \partial_x(f_2) - \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) - lf_7 = 0, \\ &\partial_t(f_4) + \partial_w(f_3) + \frac{1}{2}\partial_r(f_3) - \frac{1}{2}\partial_s(f_3) + lf_4 = 0, \\ &\partial_t(f_3) + \partial_x(f_3) + \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) + lf_4 = 0, \\ &\partial_t(f_3) + \partial_x(f_3) + \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) + lf_4 = 0, \\ &\partial_w(f_4) + \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) + \frac{1}{2}\partial_s(f_3) + lf_4 = 0, \\ &\partial_w(f_3) + \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) + \frac{1}{2}\partial_s(f_3) + lf_4 = 0, \\ &\partial_w(f_3) + \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) + \frac{1}{2}\partial_s(f_3) + lf_4 = 0, \\ &\partial_w(f_3) + \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) + \frac{1}{2}\partial_s(f_3) + lf_4 = 0, \\ &\partial_w(f_3) + \frac{1}{2}\partial_r(f_3) + \frac{1}{2}\partial_s(f_3) + \frac{1}$$

whose solution is given by

$$\begin{split} f_1(r,s,t,w,x,y,z) &= (P+R)s + (S-N)t + \frac{l}{2}\{-M(w^2+x^2) - U(y^2+z^2) + (R-P)(wy+xz) \\ &\quad + (N+S)(wz-xy) + 2Tw - 2Qx + 2Wy - 2Vz\} + C_1, \\ f_2(r,s,t,w,x,y,z) &= -(P+R)r + (M+U)t - \frac{l}{2}\{N(w^2+y^2) - S(x^2+z^2) + (R-P)(wx-yz) \\ &\quad + (M-U)(wz+xy) - 2Vw + 2Wx + 2Qy - 2Tz\} + C_2, \\ f_3(r,s,t,w,x,y,z) &= -(S-N)r - (M+U)s - \frac{l}{2}\{P(w^2+z^2) + R(x^2+y^2) + (N+S)(wx+yz) \\ &\quad + (U-M)(wy-xz) - 2Ww - 2Vx + 2Ty + 2Qz\} + C_3, \\ f_4(r,s,t,w,x,y,z) &= Mx + Ny + Pz + Q, \\ f_5(r,s,t,w,x,y,z) &= -Mw + Ry + Sz + T, \\ f_6(r,s,t,w,x,y,z) &= -Nw - Rx + Uz + V, \\ f_7(r,s,t,w,x,y,z) &= -Pw - Sx - Uy + W, \\ \end{split}$$
 where  $M,N,P,Q,R,S,T,U,V,W,C_1,C_2,C_3 \in \mathbb{R}.$ 

As a consequence, when m = 0, we obtain

Proposition 10 The Lie algebra of Killing vector fields is 13-dimensional.

## 5 Computing horizontal geodesics of the quaternionic Heisenberg group

Following the computations in [11], the vector fields

$$W = \partial_w + \frac{1}{2}(x\partial_r + y\partial_s + z\partial_t),$$

$$X = \partial_x - \frac{1}{2}(w\partial_r + z\partial_s - y\partial_t),$$

$$Y = \partial_y + \frac{1}{2}(z\partial_r - w\partial_s - x\partial_t),$$

$$Z = \partial_z - \frac{1}{2}(y\partial_r - x\partial_s - w\partial_t),$$

which are the old  $X_4, \ldots, X_7$  ones, provided m = 0, l = 1, along with  $\{\partial_r, \partial_s, \partial_t\}$ , form an orthonormal frame for the quaternionic contact manifold  $\mathbb{H} \times \text{Im}\mathbb{H}$ . This means that  $\{W, X, YZ\}$  frame the fourth plane  $\mathcal{H}$  and they are orthonormal with respect to the inner product  $ds^2 = (dw^2 + dx^2 + dy^2 + dz^2)|_{\mathcal{H}}$  on the distribution. The sub-Riemannian Hamiltonian writes down as

$$H = \frac{1}{2}(P_W^2 + P_X^2 + P_Y^2 + P_Z^2),\tag{3}$$

where  $P_W, P_X, P_Y, P_Z$  are the momentum functions of the vector fields W, X, Y, Z, respectively. Thus

$$\begin{split} P_W &= p_w + \frac{1}{2}(xp_r + yp_s + zp_t), \\ P_X &= p_x - \frac{1}{2}(wp_r + zp_s - yp_t), \\ P_Y &= p_y + \frac{1}{2}(zp_r - wp_s - xp_t), \\ P_Z &= p_z - \frac{1}{2}(yp_r - xp_s + wp_t), \end{split}$$

where  $p_w, p_x, p_y, p_z, p_r, p_s, p_t$  are the fiber coordinates on the cotangent bundle of  $\mathbb{R}^7$  corresponding to the cartesian coordinates w, x, y, z, r, s, t on  $\mathbb{R}^7$ . Again, these fiber coordinates are defined by writing a covector as  $p = p_w dw + p_x dx + p_y dy + p_z dz + p_r dr + p_s ds + p_t dt$ . Together,  $(w, x, y, z, r, s, t, p_w, p_x, p_y, p_z, p_r, p_s, p_t)$  are global coordinates on the cotangent bundle  $T^*\mathbb{R}^7 = \mathbb{R}^7 \oplus \mathbb{R}^7$ . Hamilton's equations can be written

$$\frac{df}{du} = \{f, H\}, \quad f \in C^{\infty}(T^*\mathbb{R}^7), \tag{4}$$

which holds for any smooth function f. The function H defines a vector field  $X_H$ , called the Hamiltonian vector field, which has a flow  $\Phi_u: T^*\mathbb{R}^7 \to T^*\mathbb{R}^7$ . Let  $f: T^*\mathbb{R}^7 \to \mathbb{R}$  be any smooth function on the cotangent bundle. Form the u-dependent function  $f_u = \Phi_u^* f$  by pulling f back via the flow. Thus  $f_u(w, x, y, z, r, s, t, p_w, p_x, p_y, p_z, p_r, p_s, p_t) = f(\Phi_u(w, x, y, z, r, s, t, p_w, p_x, p_y, p_z, p_r, p_s, p_t))$ . In other words  $\frac{df}{du} = X_H[f_u]$ , which gives meaning to the left-hand side of Hamilton's equations.

To define the right hand side, which is to say the vector field  $X_H$ , we will need the Poisson bracket. The Poisson bracket on the cotangent bundle  $T^*\mathbb{R}^7$  of a manifold  $\mathbb{R}^7$  is a canonical Lie algebra structure defined on the vector space  $C^{\infty}(T^*\mathbb{R}^7)$  of smooth functions on  $T^*\mathbb{R}^7$ . The Poisson bracket is denoted  $\{\cdot,\cdot\}: C^{\infty} \times C^{\infty} \to C^{\infty}$ , where  $C^{\infty} = C^{\infty}(T^*\mathbb{R}^7)$ , and can be defined by the coordinate formula

$$\{f,g\} = \frac{\partial f}{\partial x^i} \frac{\partial g}{\partial p_i} - \frac{\partial g}{\partial x^i} \frac{\partial f}{\partial p_i}.$$

This formula is valid in any canonical coordinate system, and can be shown to be coordinate independent. The Poisson bracket satisfies the Leibniz identity

$${f,gh} = g{f,h} + h{f,g},$$

which means that the operation  $\{., H\}$  defines a vector field  $X_H$ , called the Hamiltonian vector field. By letting the functions f vary over the collection of coordinate functions  $x^i$  and we get the more common form of Hamilton's equations

$$\dot{x}^{\alpha} = \frac{\partial H}{\partial p_{\alpha}}, \quad \dot{p}_{\alpha} = -\frac{\partial H}{\partial x^{\alpha}}.$$

Indeed, for the first one we take f = w and g = H. Then  $\{w, H\} = \frac{\partial w}{\partial x^i} \frac{\partial H}{\partial p_i} - \frac{\partial H}{\partial x^i} \frac{\partial w}{\partial p_i}$  if and only if  $\dot{w} = \frac{\partial H}{\partial p_m}$ . Also we have

$$\dot{x} = \frac{\partial H}{\partial p_x}, \quad \dot{y} = \frac{\partial H}{\partial p_y}, \quad \dot{z} = \frac{\partial H}{\partial p_z}.$$

These equations are in turn equivalent to the above formulation (4), which is more convenient to use, because the momentum function  $W \mapsto P_W$  is a Lie algebra anti-homomorphism from the Lie algebra of all smooth vector fields on  $\mathbb{R}^7$  to  $C(T^*\mathbb{R}^7)$  with the Poisson brackets:

$$\{P_W, P_X\} = -P_{[W,X]}, \quad \{P_W, P_Y\} = -P_{[W,Y]}, \quad \{P_W, P_Z\} = -P_{[W,Z]}, 
\{P_X, P_Y\} = -P_{[X,Y]}, \quad \{P_X, P_Z\} = -P_{[X,Z]}, \quad \{P_Y, P_Z\} = -P_{[Y,Z]}.$$
(5)

Since all calculations are similar, we only prove the first one:

$$\{P_W, P_X\} = \{p_w + \frac{x}{2}p_r + \frac{y}{2}p_s + \frac{z}{2}p_t, p_x - \frac{w}{2}p_r - \frac{z}{2}p_s + \frac{y}{2}p_t\} = p_r = -P_{[W,X]}.$$

For the quaternionic contact group, with our choose of W, X, Y, Z as a frame for  $\mathcal{H}$ , we compute

$$[W, X] = -\partial_r, \quad [W, Y] = -\partial_s, \quad [W, Z] = -\partial_t,$$

$$[X, Y] = -\partial_t, \quad [X, Z] = \partial_s, \quad [Y, Z] = -\partial_t,$$

$$[W, \partial_r] = [W, \partial_r s] = [W, \partial_t] = [X, \partial_r] = [X, \partial_s] = [X, \partial_t] = 0,$$

$$[Y, \partial_r] = [Y, \partial_s] = [Y, \partial_t] = [Z, \partial_r] = [Z, \partial_s] = [Z, \partial_t] = 0.$$

Thus

$$\{P_W, P_X\} = \partial_r := P_r, \quad \{P_W, P_Y\} = \partial_s := P_s, \quad \{P_W, P_Z\} = \partial_t := P_t,$$

$$\{P_X, P_Y\} = P_t, \quad \{P_X, P_Z\} = -p_s = -P_s, \quad \{P_Y, P_Z\} = p_r = P_r$$

We can prove that

$$\{P_W, P_r\} = \{P_W, P_s\} = \{P_W, P_t\} = \{P_X, P_r\} = \{P_X, P_s\} = \{P_X, P_t\} = 0,$$
  
$$\{P_Y, P_r\} = \{P_Y, P_s\} = \{P_Y, P_t\} = \{P_Z, P_r\} = \{P_Z, P_s\} = \{P_Z, P_t\} = 0.$$

These relations can also easily be computed by hand, from our formulae for  $P_W, P_X, P_Y, P_Z$  and the bracket in terms of  $w, x, y, z, r, s, r, p_w, p_x, p_y, p_z, p_r, p_s, p_t$ .

**Lemma 11** By letting f vary over the functions  $w, x, y, z, r, s, r, P_W, P_X, P_Y, P_Z, P_r, P_s, P_t$ , using the bracket relations and equation (5), we find that Hamilton's equations are equivalent to the system

$$\begin{array}{ll} \dot{w} = P_W, & \dot{P}_W = p_r P_X + p_s P_Y + p_t P_Z, \\ \dot{x} = P_X, & \dot{P}_X = -p_r P_W - p_s P_Z + p_t P_Y, \\ \dot{y} = P_Y, & \dot{P}_Y = p_r P_Z - p_s P_W - p_t P_X, \\ \dot{z} = P_Z, & \dot{P}_Z = -p_r P_Y + p_s P_X - p_t P_W, \\ \dot{r} = \frac{1}{2} (x P_W - w P_X + z P_Y - w P_Z), & \dot{P}_r = 0, \\ \dot{s} = \frac{1}{2} (y P_W - z P_X + x P_Y - w P_Z), & \dot{P}_s = 0, \\ \dot{t} = \frac{1}{2} (z P_W + y P_X - x P_Y - w P_Z), & \dot{P}_t = 0. \end{array}$$

To see it, remember that  $H = \frac{1}{2}(P_W^2 + P_X^2 + P_Y^2 + P_z^2)$ . Then

$$\dot{w} = \{w, H\} = P_w \frac{\partial P_W}{\partial p_w} = P_W,$$

$$\dot{x} = \{x, H\} = P_X \frac{\partial P_X}{\partial p_x} = P_X,$$

$$\dot{y} = P_Y,$$

$$\dot{z} = P_Z.$$

Also, considering that:

$$\frac{\partial P_W}{\partial p_r} = \frac{x}{2}, \qquad \frac{\partial P_W}{\partial p_s} = \frac{y}{2}, \qquad \frac{\partial P_W}{\partial p_t} = \frac{z}{2},$$

$$\frac{\partial P_X}{\partial p_r} = -\frac{w}{2}, \qquad \frac{\partial P_X}{\partial p_s} = -\frac{z}{2}, \qquad \frac{\partial P_X}{\partial p_t} = \frac{y}{2},$$

$$\frac{\partial P_Y}{\partial p_r} = \frac{z}{2}, \qquad \frac{\partial P_Y}{\partial p_s} = -\frac{w}{2}, \qquad \frac{\partial P_Y}{\partial p_t} = -\frac{x}{2},$$

$$\frac{\partial P_Z}{\partial p_r} = -\frac{y}{2}, \qquad \frac{\partial P_Z}{\partial p_s} = \frac{x}{2}, \qquad \frac{\partial P_Z}{\partial p_t} = -\frac{w}{2},$$

$$\dot{r} = \frac{1}{2}(xP_W - wP_X + zP_Y - yP_Z).$$

we have

Indeed,

$$\begin{split} \dot{r} &= \{r, H\} = P_W \frac{\partial P_W}{\partial p_r} + P_X \frac{\partial P_X}{\partial p_r} + P_Y \frac{\partial P_Y}{\partial p_r} + P_Z \frac{\partial P_Z}{\partial p_r} \\ &= \frac{1}{2} (x P_W - w P_X + z P_Y - y P_Z) \\ \dot{s} &= \{s, H\} = P_W \frac{\partial P_W}{\partial p_s} + P_X \frac{\partial P_X}{\partial p_s} + P_Y \frac{\partial P_Y}{\partial p_s} + P_Z \frac{\partial P_Z}{\partial p_s} \\ &= \frac{1}{2} (y P_W - z P_X + x P_Y - w P_Z) \\ \dot{t} &= \{t, H\} = P_W \frac{\partial P_W}{\partial p_t} + P_X \frac{\partial P_X}{\partial p_t} + P_Y \frac{\partial P_Y}{\partial p_t} + P_Z \frac{\partial P_Z}{\partial p_t} \\ &= \frac{1}{2} (z P_W + y P_X - x P_Y - w P_Z). \end{split}$$

Working as above we obtain

$$\begin{split} \dot{P}_W &= \{P_W, H\} = p_r P_X + p_s P_Y + p_t P_Z, \\ \dot{P}_X &= \{P_X, H\} = -p_r P_W - p_s P_Z + p_t P_Y, \\ \dot{P}_Y &= \{P_Y, H\} = p_r P_Z - p_s P_W - p_t P_X, \\ \dot{P}_Z &= \{P_Z, H\} = -p_r P_Y + p_s P_X - p_t P_W. \end{split}$$

Then we are ready to show the following

**Theorem 12** The horizontal geodesics of the quaternionic Heisenberg group are exactly the horizontal lifts of arcs of circles, including line segments as a degenerate case.

**Proof.** It is not difficult to see that  $\dot{P}_r = \dot{P}_s = \dot{P}_t = 0$ . These equations assert that  $P_r = p_r$ ,  $P_s = p_s$  and  $P_t = p_t$  are constant. The variables r, s, t appears nowhere in the right-hand sides of these equations. It follows that the variables  $w, x, y, z, P_W, P_X, P_Y, P_Z$  evolve independently of r, s, t, and so we can view the system as defining a one-parameter family of dynamical systems on  $\mathbb{R}^8$  parameterized by the constant value of  $P_r, P_s, P_t$ .

Combine w, x, y, z into a single quaternionic variable  $\omega = w + ix + jy + kz$  and taking into account the fourteen equations one has

$$\frac{d\omega}{du} = P_W + iP_X + jP_Y + kP_Z$$

The *u*-derivative of  $P_W + iP_X + jP_Y + kP_Z$  is  $-(ip_r + jp_s + kp_t)(P_W + iP_X + jP_Y + kP_Z)$ . Then we have  $\frac{d^2\omega}{du^2} = -(ip_r + jp_s + kp_t)\frac{d\omega}{du}$ , where  $p_r$ ,  $p_s$  and  $p_t$  are constant.

By integrating the above expression we get

$$P_W + iP_X + jP_Y + kP_Z = P(0)\exp(-(ip_r + jp_s + kp_t)t),$$

where 
$$P(0) = P_W(0) + iP_X(0) + jP_Y(0) + kP_Z(0)$$
.

A second integration yields the general form of the geodesics on the quaternionic contact group:

$$\omega(u) = w(u) + ix(u) + jy(u) + kz(u) = \frac{P(0)}{ip_r + jp_s + kp_t} \left( \exp(-(ip_r + jp_s + kp_t)t - 1) + w(0) + ix(0) + jy(0) + kz(0) \right),$$

$$r(u) = r(0) + \frac{1}{2} \int_0^t \operatorname{Im}_I(\bar{\omega} \, d\omega),$$

$$s(u) = s(0) + \frac{1}{2} \int_0^t \operatorname{Im}_J(\bar{\omega} \, d\omega),$$

$$r(u) = t(0) + \frac{1}{2} \int_0^t \operatorname{Im}_K(\bar{\omega} \, d\omega).$$

## 6 Appendix

The brackets:

$$\begin{split} [X_4,X_5] &= -l\{1+m(y^2+z^2)\}X_1+ml(wz+xy)X_2-ml(wy-xz)X_3-2mxX_4+2mwX_5,\\ [X_4,X_6] &= -ml(wz-xy)X_1-l\{1+m(x^2+z^2)\}X_2+ml(wx+yz)X_3-2myX_4+2mwX_6,\\ [X_4,X_7] &= ml(wy+xz)X_1-ml(wx-yz)X_2-l\{1+(x^2+y^2)\}X_3-2mzX_4+2mwX_7,\\ [X_5,X_6] &= -ml(wy+xz)X_1+ml(wx-yz)X_2-l\{1+m(w^2+z^2)\}X_3-2myX_5+2mxX_6,\\ [X_5,X_7] &= ml(xy-wz)X_1+l\{1+m(w^2+y^2)\}X_2+ml(wx+yz)X_3-2mzX_5+2mxX_7,\\ [X_6,X_7] &= -l\{1+m(w^2+x^2)\}X_1-ml(wz+xy)X_2+ml(wy-xz)X_3-2mzX_6+2myX_7. \end{split}$$

The Levi-Civita connection:

$$\begin{split} &\nabla_{X_1}X_4 = \frac{1}{2}\{1 + m(y^2 + z^2)\}X_5 + \frac{ml}{2}(wz - xy)X_6 - \frac{ml}{2}(wy + xz)X_7, \\ &\nabla_{X_1}X_5 = -\frac{1}{2}\{1 + m(y^2 + z^2)\}X_4 + \frac{ml}{2}(wy + xz)X_5 + \frac{ml}{2}(wz - xy)X_7, \\ &\nabla_{X_1}X_6 = -\frac{ml}{2}(wz - xy)X_4 - \frac{ml}{2}(wy + xz)X_5 + \frac{1}{2}\{1 + m(w^2 + x^2)\}X_7, \\ &\nabla_{X_1}X_7 = \frac{ml}{2}(wy + xz)X_4 - \frac{ml}{2}(wz - xy)X_5 - \frac{1}{2}\{1 + m(w^2 + x^2)\}X_6, \\ &\nabla_{X_2}X_4 = -\frac{ml}{2}(wz + xy)X_5 + \frac{1}{2}\{1 + m(x^2 + z^2)\}X_6 + \frac{ml}{2}(wx - yz)X_7, \\ &\nabla_{X_2}X_5 = \frac{ml}{2}(wz + xy)X_4 - \frac{ml}{2}(wx - yz)X_6 - \frac{1}{2}\{1 + m(w^2 + y^2)\}X_7, \\ &\nabla_{X_2}X_5 = -\frac{1}{2}\{1 + m(x^2 + z^2)\}X_4 + \frac{ml}{2}(wx - yz)X_5 + \frac{ml}{2}(wz + xy)X_7, \\ &\nabla_{X_2}X_7 = -\frac{ml}{2}(wx - yz)X_4 + \frac{1}{2}\{1 + m(w^2 + y^2)\}X_5 - \frac{ml}{2}(wz + xy)X_6, \\ &\nabla_{X_3}X_4 = \frac{ml}{2}(wy - xz)X_5 - \frac{ml}{2}(wx + yz)X_6 + \frac{1}{2}\{1 + m(x^2 + y^2)\}X_7, \\ &\nabla_{X_3}X_5 = -\frac{ml}{2}(wy - xz)X_4 + \frac{1}{2}\{1 + m(w^2 + z^2)\}X_6 - \frac{ml}{2}(wx + yz)X_7, \\ &\nabla_{X_3}X_6 = \frac{ml}{2}(wx + yz)X_4 - \frac{1}{2}\{1 + m(w^2 + z^2)\}X_5 - \frac{ml}{2}(wy - xz)X_7, \\ &\nabla_{X_3}X_7 = -\frac{1}{2}\{1 + m(x^2 + y^2)\}X_4 + \frac{ml}{2}(wx + yz)X_5 + \frac{ml}{2}(wy - xz)X_6, \\ &\nabla_{X_4}X_4 = 2m(xX_5 + yX_6 + zX_7), \\ &\nabla_{X_4}X_5 = -\frac{1}{2}\{1 + m(y^2 + z^2)\}X_1 + \frac{ml}{2}(wz + xy)X_2 - \frac{ml}{2}(wx + yz)X_3 - 2mxX_4, \\ &\nabla_{X_4}X_7 = \frac{ml}{2}(wz - xy)X_1 - \frac{1}{2}\{1 + m(x^2 + z^2)\}X_2 + \frac{ml}{2}(wx + yz)X_3 - 2mxX_4, \\ &\nabla_{X_5}X_4 = \frac{1}{2}\{1 + m(y^2 + z^2)\}X_1 - \frac{ml}{2}(wz - yz)X_2 - \frac{1}{2}\{1 + m(x^2 + y^2)\}X_3 - 2mzX_4, \\ &\nabla_{X_5}X_6 = -\frac{ml}{2}(wy + xz)X_1 - \frac{ml}{2}(wz - yz)X_2 - \frac{1}{2}\{1 + m(w^2 + z^2)\}X_3 - 2myX_5, \\ &\nabla_{X_5}X_6 = -\frac{ml}{2}(wy + xz)X_1 + \frac{ml}{2}(wx - yz)X_2 - \frac{1}{2}\{1 + m(w^2 + z^2)\}X_3 - 2myX_5, \\ &\nabla_{X_5}X_6 = -\frac{ml}{2}(wy + xz)X_1 + \frac{ml}{2}(wx - yz)X_2 - \frac{1}{2}\{1 + m(w^2 + z^2)\}X_3 - 2myX_5, \\ &\nabla_{X_5}X_6 = -\frac{ml}{2}(wy + xz)X_1 + \frac{ml}{2}(wx - yz)X_2 - \frac{1}{2}\{1 + m(w^2 + z^2)\}X_3 - 2myX_5, \\ &\nabla_{X_5}X_6 = -\frac{ml}{2}(wy + xz)X_1 + \frac{ml}{2}(wx - yz)X_2 - \frac{1}{2}\{1 + m(w^2 + z^2)\}X_3 - 2myX_5, \\ &\nabla_{X_5}X_7 = -\frac{ml}{2}(wz - xy)X_1 + \frac{1}{2}\{1 + m(w^2 + y^2)X_2 - \frac$$

$$\nabla_{X_{6}}X_{4} = \frac{ml}{2}(wz - xy)X_{1} + \frac{l}{2}\{1 + m(x^{2} + z^{2})\}X_{2} - \frac{ml}{2}(wx + yz)X_{3} - 2mwX_{6},$$

$$\nabla_{X_{6}}X_{5} = \frac{ml}{2}(wy + xz)X_{1} - \frac{ml}{2}(wx - yz)X_{2} + \frac{l}{2}\{1 + m(w^{2} + z^{2})\}X_{3} - 2mxX_{6},$$

$$\nabla_{X_{6}}X_{6} = 2m(wX_{4} + xX_{5} + zX_{7}),$$

$$\nabla_{X_{6}}X_{7} = -\frac{l}{2}\{1 + m(w^{2} + x^{2})\}X_{1} - \frac{ml}{2}(wz + xy)X_{2} + \frac{ml}{2}(wy - xz)X_{3} - 2mzX_{6},$$

$$\nabla_{X_{7}}X_{4} = -\frac{ml}{2}(wy + xz)X_{1} + \frac{ml}{2}(wx - yz)X_{2} + \frac{l}{2}\{1 + m(x^{2} + y^{2})\}X_{3} - 2mwX_{7},$$

$$\nabla_{X_{7}}X_{5} = \frac{ml}{2}(wz - xy)X_{1} - \frac{l}{2}\{1 + m(w^{2} + y^{2})\}X_{2} - \frac{ml}{2}(wx + yz)X_{3} - 2mxX_{7},$$

$$\nabla_{X_{7}}X_{6} = \frac{l}{2}\{1 + m(w^{2} + x^{2})\}X_{1} + \frac{ml}{2}(wz + xy)X_{2} - \frac{ml}{2}(wy - xz)X_{3} - 2myX_{7},$$

$$\nabla_{X_{7}}X_{7} = 2m(wX_{4} + xX_{5} + yX_{6}).$$

The curvature tensor:

$$R_{X_1X_4X_1X_4} = R_{X_1X_5X_1X_5} = \frac{l^2}{4} \{1 + m(K+1)(y^2 + z^2)\},$$

$$R_{X_1X_6X_1X_6} = R_{X_1X_7X_1X_7} = \frac{l^2}{4} \{1 + m(K+1)(w^2 + x^2)\},$$

$$R_{X_2X_4X_2X_4} = R_{X_2X_6X_2X_6} = \frac{l^2}{4} \{1 + m(K+1)(x^2 + z^2)\},$$

$$R_{X_2X_5X_2X_5} = R_{X_2X_7X_2X_7} = \frac{l^2}{4} \{1 + m(K+1)(w^2 + y^2)\},$$

$$R_{X_3X_4X_3X_4} = R_{X_3X_7X_3X_7} = \frac{l^2}{4} \{1 + m(K+1)(x^2 + y^2)\},$$

$$R_{X_3X_5X_3X_5} = R_{X_3X_6X_3X_6} = \frac{l^2}{4} \{1 + m(K+1)(w^2 + z^2)\},$$

$$R_{X_4X_5X_4X_5} = 4m - 3R_{X_1X_4X_1X_4},$$

$$R_{X_4X_5X_4X_5} = 4m - 3R_{X_2X_4X_2X_4},$$

$$R_{X_4X_7X_4X_7} = 4m - 3R_{X_3X_5X_3X_5},$$

$$R_{X_5X_7X_5X_7} = 4m - 3R_{X_2X_5X_2X_5},$$

$$R_{X_6X_7X_6X_7} = 4m - 3R_{X_1X_6X_1X_6}.$$

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