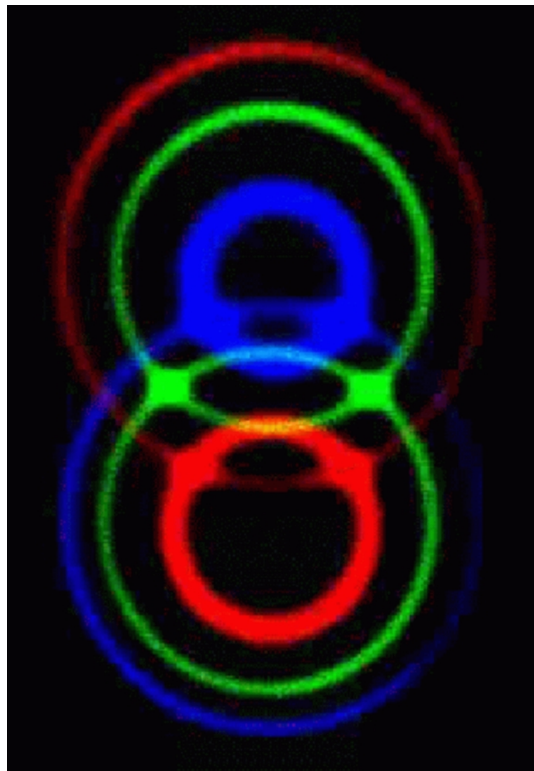


# Introduction to Quantum Chemistry

TEMA: INTRODUCTION TO QUANTUM CHEMISTRY

Introduction



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# I.A. Discrete spectrum: Definition of probability



- Discrete spectrum  $\Rightarrow$  Finite number of possible measurement results

Ex. Coin toss: Spectrum = 1,2

Ex. Dice toss: Spectrum = 1,2,3,4,5,6

- Probability  $P_N \Rightarrow$  Frequency of every result

$$P_N = \frac{\text{Number of measurements providing } N \text{ as result}}{\text{Number of possible measurements}}$$

- Normalization  $\Rightarrow \sum_i P_i = 1$

Ex. Coin toss:  $P_1 = \frac{1}{2} = P_2$

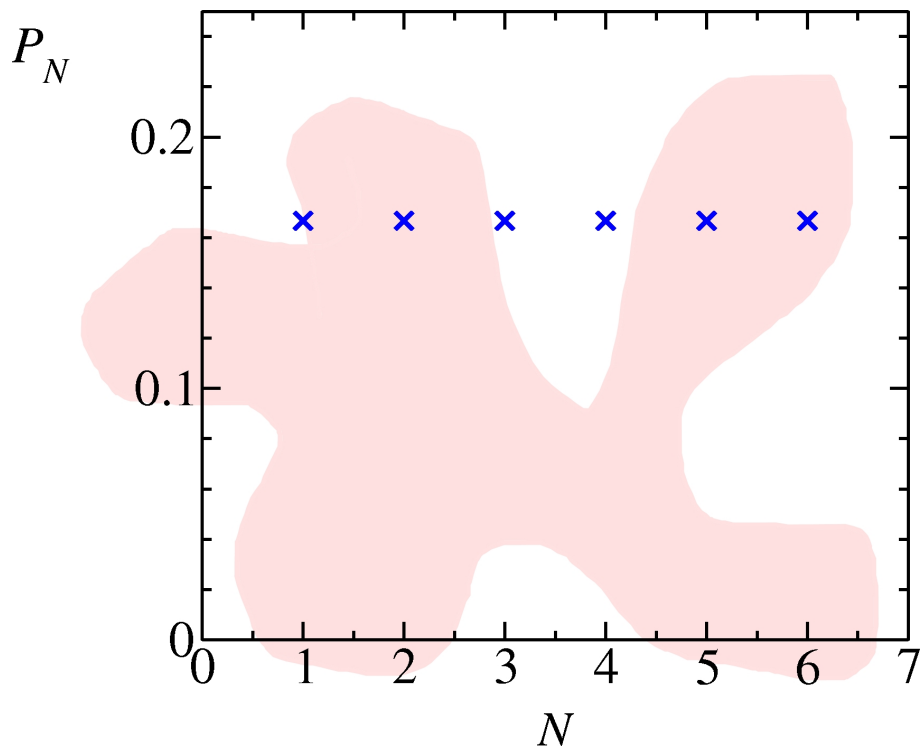
Ex. Dice toss:  $P_1 = \frac{1}{6} = P_2 = P_3 = P_4 = P_5 = P_6$



# I.A. Discrete spectrum: Graphical representation



Ex. Dice toss:  $P_1 = P_2 = P_3 = P_4 = P_5 = P_6 = \frac{1}{6}$

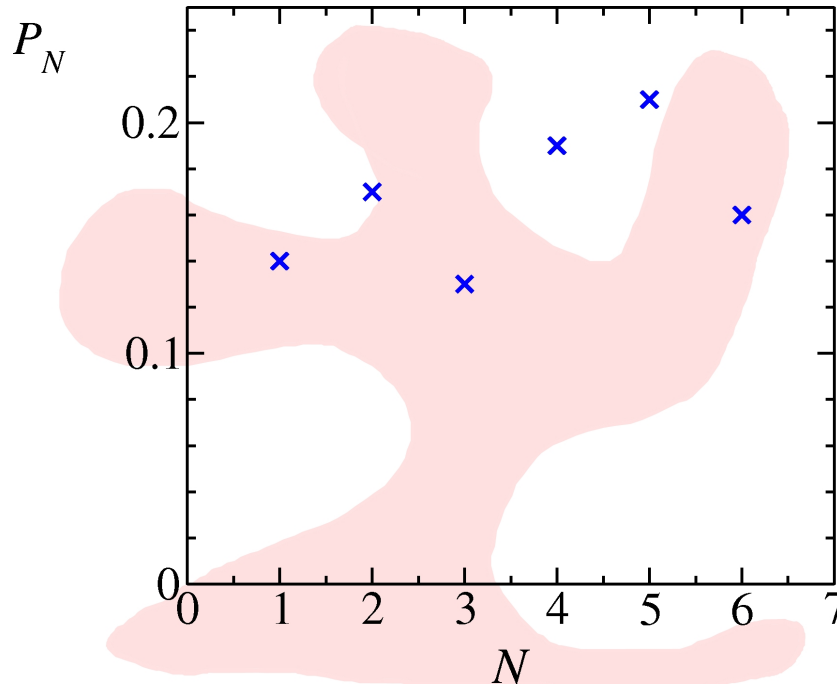




# I.A. Discrete spectrum: Graphical representation

Ex. Real dice toss:

$$P_1 = 0.14, P_2 = 0.17, P_3 = 0.13, P_4 = 0.19, P_5 = 0.21, P_6 = 0.16$$





# I.A. Discrete spectrum: Mean value of a measurement

- Mean value of a measurement  $\langle f \rangle$

N measurements  $\Rightarrow f_5, f_3, f_1, f_1, f_3, \dots$

$$\begin{aligned}\text{Mean value} &= \frac{1}{N} (f_5 + f_3 + f_1 + f_1 + f_3 + \dots) \\ &= \frac{1}{N} (f_1 N_1 + f_2 N_2 + \dots) \\ &= f_1 \frac{N_1}{N} + f_2 \frac{N_2}{N} + \dots\end{aligned}$$

$$\langle f \rangle = \sum_i f_i P_i$$



# I.A. Discrete spectrum: Mean value of a measurement



Ex. Dice toss:

$$\langle N \rangle = 1 \cdot \frac{1}{6} + 2 \cdot \frac{1}{6} + 3 \cdot \frac{1}{6} + 4 \cdot \frac{1}{6} + 5 \cdot \frac{1}{6} + 6 \cdot \frac{1}{6} = \frac{21}{6} = 3.5$$

$$\begin{aligned} \langle \sqrt{N} \rangle &= \sqrt{1} \cdot \frac{1}{6} + \sqrt{2} \cdot \frac{1}{6} + \sqrt{3} \cdot \frac{1}{6} + \sqrt{4} \cdot \frac{1}{6} + \sqrt{5} \cdot \frac{1}{6} + \sqrt{6} \cdot \frac{1}{6} \\ &= 1.8 \end{aligned}$$

$$\sqrt{\langle N \rangle} \neq \langle \sqrt{N} \rangle$$

$$\sqrt{3.5} \neq 1.8$$



# I.A. Discrete spectrum: Mean squared deviation

- Mean squared deviation  $\Delta f \Rightarrow$  Measures the mean deviation of the values from their mean value

Measurements:  $f_1, f_2, f_3, \dots$

Mean value of the deviations =  $\langle f - \langle f \rangle \rangle = \langle f \rangle - \langle f \rangle = 0$

$$\begin{aligned}(\Delta f)^2 &= \langle (f - \langle f \rangle)^2 \rangle = \langle f^2 - 2f\langle f \rangle + \langle f \rangle^2 \rangle \\ &= \langle f^2 \rangle - 2\langle f \rangle\langle f \rangle + \langle f \rangle^2\end{aligned}$$

$$\Delta f = \sqrt{\langle f^2 \rangle - \langle f \rangle^2}$$



# I.B. Continuous spectrum: Probability density

- Continuous spectrum  $\Rightarrow$  Infinite number of possible results of the measurements  $\tau \in (a, b)$ . It makes no sense to ask about the probability of a particular outcome.
- Probability density  $\rho_\tau \Rightarrow$  It describes how the probability is distributed among the possible results of the measurement

$$\rho_\tau = \frac{dP}{d\tau}$$

$$P(\tau \in (\tau_1, \tau_2)) = \int_{\tau_1}^{\tau_2} \rho_\tau d\tau$$

- Normalization  $\Rightarrow \int_{\forall \tau} \rho_\tau d\tau = 1$

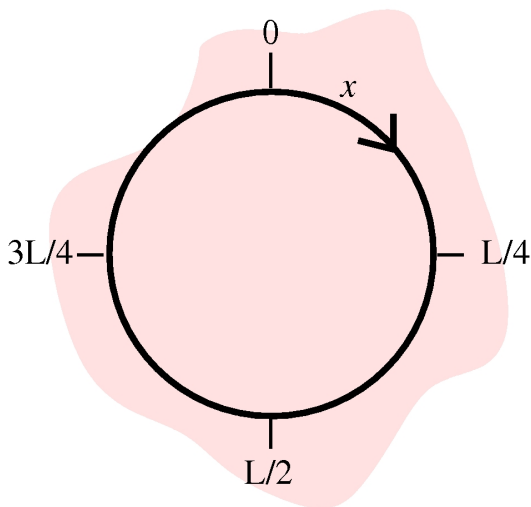




# I.B. Continuous spectrum: Probability density



Ex. Ring toss  $\Rightarrow \rho_x = \frac{dP}{dx}$



$$P(x \in [x_1, x_2]) = \int_{x_1}^{x_2} \rho_x dx$$

$$\text{Normalization} \Rightarrow \int_0^L \rho_x dx = 1$$

$$\rho_x = \text{cte.} \Rightarrow \rho_x \int_0^L dx = 1 \Rightarrow \rho_x = \frac{1}{L}$$

$$\begin{aligned} P(x \in [0, L/2]) &= \int_0^{L/2} \rho_x dx = \frac{1}{L} \int_0^{L/2} dx \\ &= \frac{1}{2} \end{aligned}$$



# I.B. Continuous spectrum: Mean value of measurement

- Mean value of a measurement  $\langle f \rangle$

$$\langle f \rangle = \int_{\forall \tau} f_{\tau} \rho_{\tau} d\tau$$

Ex. Ring toss

$$\langle x \rangle = \int_0^L x \rho_x dx = \frac{1}{L} \int_0^L x dx = \frac{L}{2}$$

$$\langle x^2 \rangle = \int_0^L x^2 \rho_x dx = \frac{1}{L} \int_0^L x^2 dx = \frac{L^2}{3}$$



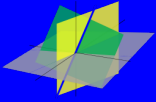
# I.B. Continuous spectrum: Mean squared deviation

- Mean squared deviation  $\Delta f$

$$\Delta f = \sqrt{\langle f^2 \rangle - \langle f \rangle^2}$$

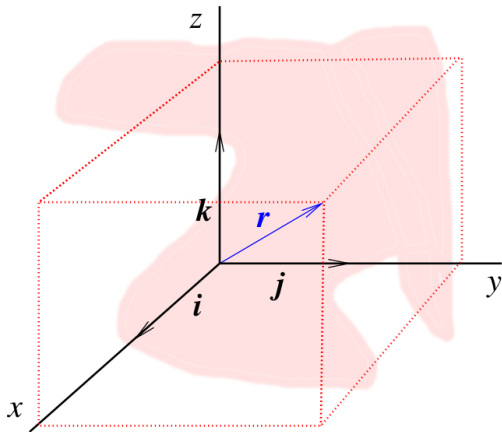
Ex. Ring toss

$$\Delta x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \frac{L}{\sqrt{12}}$$



# II.A. Vectorial space

- Vector  $\Rightarrow$  Member of a vectorial space



base  $(i, j, k)$

Normalized

$$|i|^2 = i \cdot i = 1 \text{ (idem } j, k)$$

Orthogonal/Perpendicular

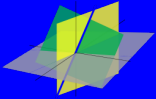
$$i \cdot j = i \cdot k = \dots = 0$$

Generator system

$$r = ai + bj + ck$$

$$\left. \begin{aligned} r \cdot i &= a \cancel{i \cdot i} + b \cancel{j \cdot i} + c \cancel{k \cdot i} = a \\ r \cdot j &= a \cancel{i \cdot j} + b \cancel{j \cdot j} + c \cancel{k \cdot j} = b \\ r \cdot k &= a \cancel{i \cdot k} + b \cancel{j \cdot k} + c \cancel{k \cdot k} = c \end{aligned} \right\}$$

$$r = (r \cdot i)i + (r \cdot j)j + (r \cdot k)k$$



# II.A. Vectorial space

- Vector  $\Rightarrow f(x) \Rightarrow$  Hilbert's space

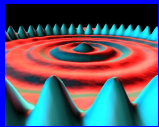
Dot product  $\Rightarrow \langle f|g \rangle = \int_{\forall x} f^*(x)g(x)dx$



David Hilbert

base  $(\{\phi_i(x)\}_{i=1}^N)$   $\left\{ \begin{array}{l} \text{Normalized} \\ \langle \phi_i | \phi_i \rangle = 1, \quad i = 1, \dots, N \\ \text{Orthogonal/Perpendicular} \\ \langle \phi_i | \phi_j \rangle = 0, \quad i \neq j = 1, \dots, N \\ \text{Generator system} \\ f(x) = \sum_{i=1}^N a_i \phi_i(x) \end{array} \right.$

$$\left. \langle \phi_j | f \rangle = \sum_{i=1}^N a_i \langle \phi_j | \phi_i \rangle = a_j \right\} f(x) = \sum_{i=1}^N \langle \phi_i | f \rangle \phi_i(x)$$



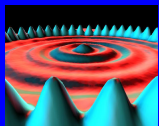
## II.B. Concept of state: Classic state

- Classical Mechanics  $\Rightarrow$  The state of the system is characterized by the positions and moments of all the particles in the system  $(\vec{r}, \vec{p})$

$$\vec{p} = m \frac{d\vec{r}}{dt}$$

If these magnitudes are known at a given time  $(\vec{r}(0), \vec{p}(0))$  they can be known at any later  $(\vec{r}(t), \vec{p}(t))$  or earlier time  $(\vec{r}(-t), \vec{p}(-t))$  using the second Newton's law

$$\vec{F} = m \frac{d^2\vec{r}}{dt^2}$$



## II.B. Concept of state: Quantum state

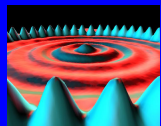
- Quantum Mechanics  $\Rightarrow$  Postulate 1. The state of the system is characterized by a function depending on the coordinates of the particles in the system and time  $\psi(\vec{r}, t)$  referred to as state function.

The squared modulus of the state function  $|\psi(\vec{r}, t)|^2$  is the probability density of the system

$$dP(\vec{r} \in [\vec{r}, \vec{r} + d\vec{r}]) = |\psi(\vec{r}, t)|^2 d\vec{r}$$

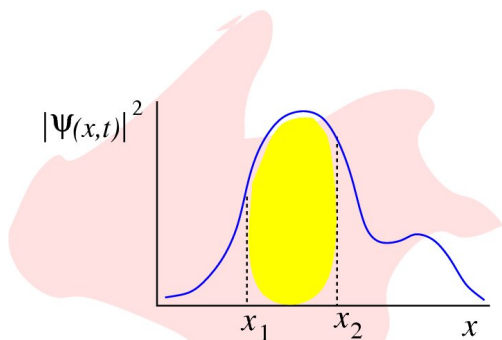
Modulus of a complex number

$$a = a_r + ia_i \Rightarrow |a|^2 = a \cdot a^* = (a_r + ia_i)(a_r - ia_i) = a_r^2 + a_i^2$$



## II.C. State function

- Probabilistic interpretation



$$P(x \in [x_1, x_2]) = \int_{x_1}^{x_2} |\Psi(x, t)|^2 dx$$

- State function must be *well behaved*:

a) Normalizable  $\Rightarrow \int_{\forall \vec{r}} |\Psi(\vec{r}, t)|^2 d\vec{r} = 1$

b) Single valued  $\Rightarrow$  Probability density has only one value in any space point

c) Continuous  $\Rightarrow$  No jumps





## II.D. Operators

- Postulate 2. Every measurable physical quantity is described by a **linear and hermitian operator** obtained from the classic expression of the magnitude using the **correspondence principle**
  - $x \rightarrow \hat{x}$
  - $p_x \rightarrow \hat{p}_x = \frac{\hbar}{i} \frac{d}{dx}$



## II.D. Operators

- Operator

$$\begin{array}{ccc} \hat{A} = \frac{d}{dx} & f(x) = e^{-2x} & \hat{B} = x^2 \\ \downarrow & & \downarrow \\ \hat{A}f = \frac{d}{dx}e^{-2x} = -2e^{-2x} & & \hat{B}f = x^2e^{-2x} \end{array}$$

- Lineal  $\begin{cases} \hat{A}(f(x) + g(x)) = \hat{A}f(x) + \hat{A}g(x) \\ \hat{A}(cf(x)) = c\hat{A}f(x) \end{cases}$

- Hermitian  $\Rightarrow \langle f|\hat{A}g\rangle = \langle \hat{A}f|g\rangle$

$$\langle f|\hat{A}g\rangle = \int_{\forall x} f^* \hat{A}g dx$$

$$\langle \hat{A}f|g\rangle = \int_{\forall x} (\hat{A}f)^* g dx = \left[ \int_{\forall x} (\hat{A}f) g^* dx \right]^* = \left[ \int_{\forall x} g^* (\hat{A}f) dx \right]^* = \langle g|\hat{A}f\rangle^*$$



## II.D. Operators

Ex. Kinetic energy operator

$$T = \frac{p^2}{2m} \rightarrow \hat{T} = \frac{\hat{p}^2}{2m} = \frac{1}{2m} \left( \frac{\hbar}{i} \frac{d}{dx} \right)^2 = \frac{-\hbar^2}{2m} \frac{d^2}{dx^2}$$

Ex. Potential energy operator

$$V(x) \rightarrow \hat{V} = V(x)$$

Ex. Hamiltonian operator

$$E = \frac{p^2}{2m} + V(x) \rightarrow \hat{H} = \hat{T} + \hat{V}(x) = \frac{-\hbar^2}{2m} \frac{d^2}{dx^2} + \hat{V}(x)$$



## II.D. Operators: Result of a measurement

- Postulado 3. The only possible result of the measurement of a physical quantity  $A$  is one of the eigenvalues of the corresponding eigenvalue equation

$$\hat{A} \varphi_i = a_i \varphi_i \quad i = 1, 2, \dots$$

The numbers  $a_1, a_2, \dots$  are known as eigenvalues of the  $\hat{A}$  operator and the functions  $\varphi_i$  are their corresponding eigenfunctions. The eigenfunctions define a basis set in the Hilbert's space.

$$\psi = \sum_i \langle \varphi_i | \psi \rangle \varphi_i$$



## II.D. Operators: Result of a measurement

- When the system is described by the state function  $\psi$ , the measurement of the  $A$  magnitude will provide the result  $a_i$  with a probability equal to  $P_{a_i} = |\langle \phi_i | \psi \rangle|^2$ .

$$\psi = \sum_i \langle \phi_i | \psi \rangle \phi_i \Rightarrow P_{a_i} = |\langle \phi_i | \psi \rangle|^2$$



## II.D. Operators: Result of a measurement

- Consequence 1  $\Rightarrow \hat{A}\varphi_i = a_i\varphi_i$  y  $\Psi = \varphi_i$

$$P_{a_i} = |\langle \varphi_i | \Psi \rangle|^2 = 1$$

- Consequence 2  $\Rightarrow$  the mean value of a measurement of the  $A$  magnitude is

$$\langle \hat{A} \rangle = \sum_i P_{a_i} a_i = \sum_i |\langle \varphi_i | \Psi \rangle|^2 a_i = \langle \Psi | \hat{A} | \Psi \rangle$$

$$\Rightarrow \text{Mean squared deviation } \Delta A = \sqrt{\langle \hat{A}^2 \rangle - \langle \hat{A} \rangle^2}$$



## II.D. Operators

- Consequence 3  $\Rightarrow f \equiv g, \hat{A}f = af$

$$\langle f|\hat{A}f\rangle = [\langle f|\hat{A}f\rangle]^*$$

$$\langle f|af\rangle = [\langle f|af\rangle]^*$$

$$a\langle f|f\rangle = a^*[\langle f|f\rangle]^* \Rightarrow a = a^* \Rightarrow a \text{ real}$$

- Consequence 4  $\Rightarrow \hat{A}f = af, \hat{A}g = bg, a \neq b$

$$\langle f|\hat{A}g\rangle = [\langle g|\hat{A}f\rangle]^*$$

$$b\langle f|g\rangle = a\langle g|f\rangle^*$$

$$\downarrow \langle g|f\rangle^* = [\int_{\forall x} g^* f dx]^* = \int_{\forall x} f^* g dx = \langle f|g\rangle$$

$$(b - a)\langle f|g\rangle = 0 \Rightarrow f, g \text{ orthogonals}$$



## II.D. Operators: Commutators

- The commutator between two operators is defined by

$$[\hat{A}, \hat{B}] = \hat{A}\hat{B} - \hat{B}\hat{A}$$

Ex.  $\hat{A} = \frac{d}{dx}$  y  $\hat{B} = x^2$

$$\begin{aligned} [\hat{A}, \hat{B}]f &= \hat{A}\hat{B}f - \hat{B}\hat{A}f = \frac{d}{dx} (x^2 f) - x^2 \frac{df}{dx} \\ &= 2xf + \cancel{x^2 \frac{df}{dx}} - \cancel{x^2 \frac{df}{dx}} = 2xf \end{aligned}$$

$$[\hat{A}, \hat{B}] = 2x$$





## II.D. Operators: Simultaneous measurement

- **Consequence.** If the hermitian operators  $\hat{A}$  and  $\hat{B}$  commute, then there is a common set of eigenfunctions.

$$[\hat{A}, \hat{B}] = 0 \Rightarrow \hat{A}\hat{B}f = \hat{B}\hat{A}f$$

$$\hat{A}f_i = a_i f_i$$

↓

$$\hat{B}\hat{A}f_i = \hat{B}(a_i f_i)$$

$$\hat{A}(\hat{B}f_i) = a_i(\hat{B}f_i)$$

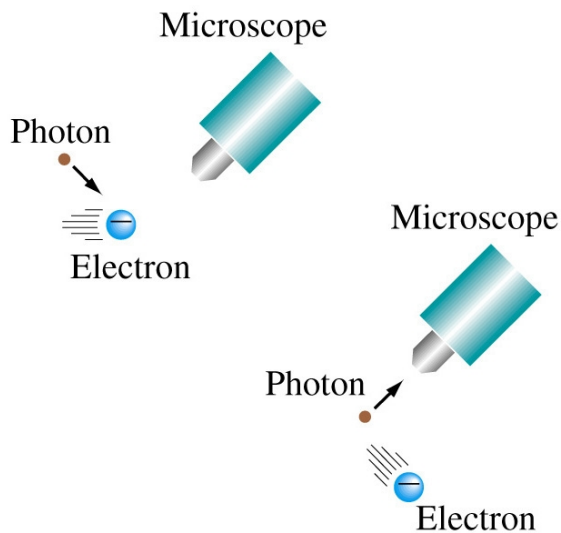
↓  $\hat{B}f_i$  is eigenfunction of  $\hat{A}$  with eigenvalue  $a_i$

$$\hat{B}f_i = k f_i$$



## II.E. Uncertainty principle

- Heisenberg's uncertainty principle



$$\Delta A \Delta B \geq \frac{1}{2} \langle [\hat{A}, \hat{B}] \rangle$$

$$[\hat{x}, \hat{p}] = i\hbar$$

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$