



# Useful quantum models in Chemistry

Adolfo Bastida



PHYSICAL CHEMISTRY I

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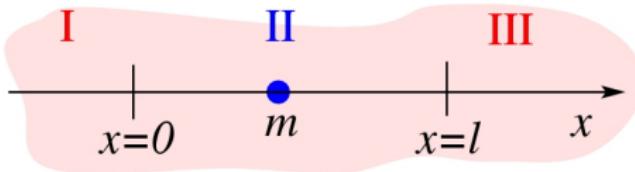
- Model
- Eigenfunctions and eigenvalues
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## 3 Particle on a sphere

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# Schrödinger's equation



$$V(x) = \begin{cases} +\infty & x \leq 0 \\ 0 & 0 < x < L \\ +\infty & L \leq x \end{cases}$$

(I)      (II)      (III)

- Zones (I) and (III)  $\Rightarrow |\psi(x)|^2 = 0 \rightarrow \psi_I(x) = \psi_{III}(x) = 0$
  - Zone (II)

$$\hat{H}\psi = E\psi$$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + \underbrace{V(x)\psi(x)}_0 = E\psi(x)$$



# Schrödinger's equation

$$\frac{d^2\psi(x)}{dx^2} = -\frac{2mE}{\hbar^2}\psi(x)$$

- Independent solutions

$$\psi_1(x) = \sin(kx)$$

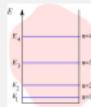
$$\psi_2(x) = \cos(kx)$$

$$\frac{d^2\psi_1(x)}{dx^2} = -k^2 \sin(kx)$$

$$\frac{d^2\psi_2(x)}{dx^2} = -k^2 \cos(kx)$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$

$$\psi(x) = A\psi_1(x) + B\psi_2(x) = A \sin\left(\frac{\sqrt{2mE}}{\hbar}x\right) + B \cos\left(\frac{\sqrt{2mE}}{\hbar}x\right)$$



# Eigenfunctions and eigenvalues

$$\psi(x) = A \sin\left(\frac{\sqrt{2mE}}{\hbar}x\right) + B \cos\left(\frac{\sqrt{2mE}}{\hbar}x\right)$$

- Continuity  $x = 0 \Rightarrow \psi_{\text{I}}(x = 0) = \psi_{\text{II}}(x = 0)$

$$A \sin\left(\frac{\sqrt{2mE}}{\hbar}0\right) + B \cos\left(\frac{\sqrt{2mE}}{\hbar}0\right) = 0$$

$$A \cdot 0 + B \cdot 1 = 0$$

$$B = 0$$

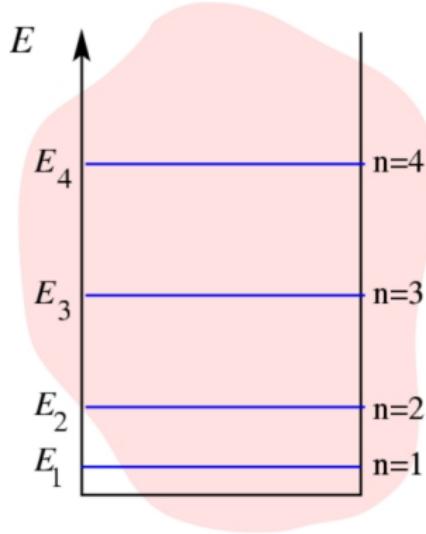
- Continuity  $x = L \Rightarrow \psi_{\text{II}}(x = L) = \psi_{\text{III}}(x = L)$

$$A \cdot \sin\left(\frac{\sqrt{2mE}}{\hbar}L\right) = 0 \Rightarrow \frac{\sqrt{2mE}}{\hbar}L = n\pi; \quad n = 0, \pm 1, \pm 2, \dots$$

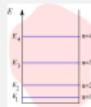


## Quantization

$$\psi_{\text{II}}(x) = A \sin\left(\frac{n\pi x}{L}\right); \quad n = 1, 2, \dots \quad \begin{cases} n = 0 & \text{no normalizable} \\ n < 0 & \text{phase factor} \end{cases}$$



- Quantization  $E_n = \frac{\hbar^2}{8mL^2} n^2$   
 $n = 1, 2, \dots$
- Ground state  $\Rightarrow E_1$
- Spacing  $\Rightarrow E_n = n^2 E_1$



# Normalization

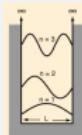
- Normalization

$$\int_{-\infty}^{+\infty} |\psi(x)|^2 dx = 1$$

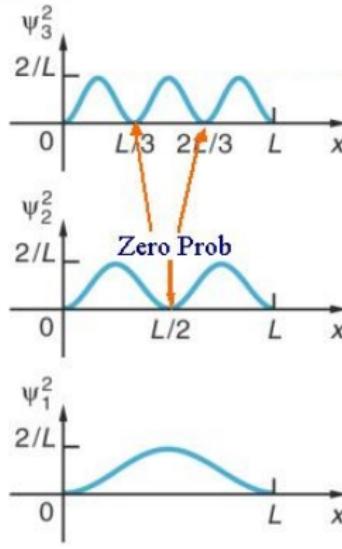
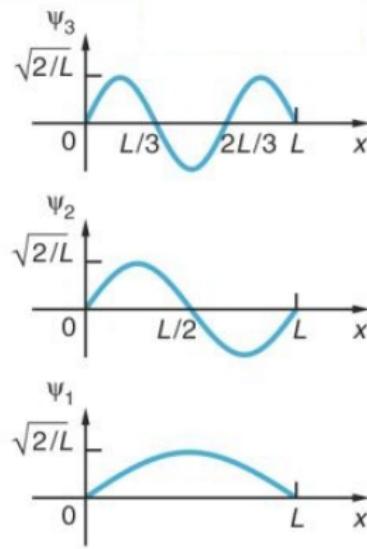
$$\underbrace{\int_{-\infty}^0 |\psi_I(x)|^2 dx}_{0} + \int_0^L |\psi_{II}(x)|^2 dx + \underbrace{\int_L^{+\infty} |\psi_{III}(x)|^2 dx}_{0} = 1$$

$$|A|^2 \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx = 1 \Rightarrow |A| = \sqrt{\frac{2}{L}}$$

$$\psi_{II}(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$



## Plots



- Nodes  $\psi(x_i) = 0$



# Classical limit

Ex. electron vs melon

- Electron in a box

$$m = 9.1 \cdot 10^{-31} \text{ kg}$$

$$L = 10 \text{ \AA}$$

$$E_n = 6 \cdot 10^{-20} n^2 \text{ J}$$

$$\Delta E \sim 10^{-20} \text{ J}$$

- Melon in a box

$$m = 1 \text{ kg}$$

$$L = 1 \text{ m}$$

$$E_n = 4 \cdot 10^{-68} n^2 \text{ J}$$

$$\Delta E \sim 10^{-68} \text{ J}$$

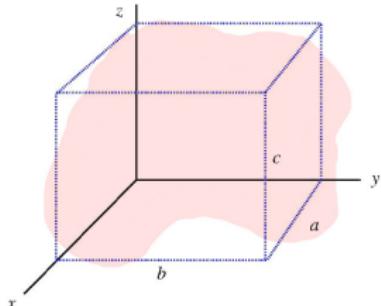
$$v = 1 \text{ m/s} \Rightarrow E = 0.5 \text{ J}$$

$$n = \sqrt{\frac{8mL^2E}{h^2}} \sim 10^{33}$$



## 3D box

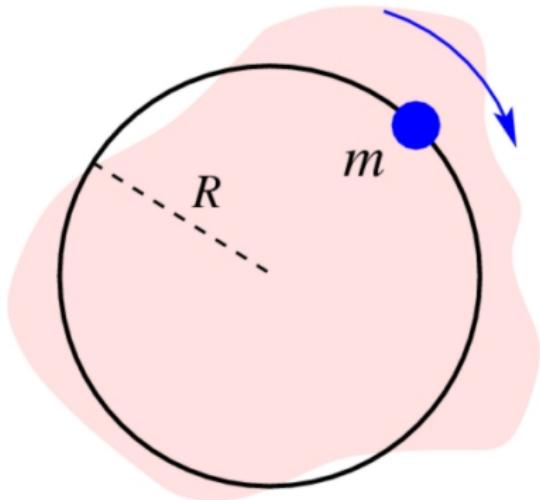
$$\hat{H}_{3D}(x, y, z) = -\frac{\hbar^2}{2m} \left( \frac{d^2}{dx^2} + \frac{d^2}{dy^2} + \frac{d^2}{dz^2} \right) \begin{cases} 0 \leq x \leq a \\ 0 \leq y \leq b \\ 0 \leq z \leq c \end{cases}$$



- $\hat{H}_{3D}\Psi_{3D} = E_{3D}\Psi_{3D}$
- $\hat{H}_{3D}(x, y, z) = \hat{H}_{1D}(x) + \hat{H}_{1D}(y) + \hat{H}_{1D}(z)$
- $\Psi_{3D}(x, y, z) = \Psi_{1D}(x)\Psi_{1D}(y)\Psi_{1D}(z)$
- $\hat{H}_{1D}(\alpha)\Psi_{1D}(\alpha) = E_{1D}\Psi_{1D}(\alpha)$  ( $\alpha = x, y, z$ )
- $E_{3D} = \frac{n_x^2 h^2}{8ma^2} + \frac{n_y^2 h^2}{8mb^2} + \frac{n_z^2 h^2}{8mc^2}$      $n_x, n_y, n_z = 1, 2, \dots$



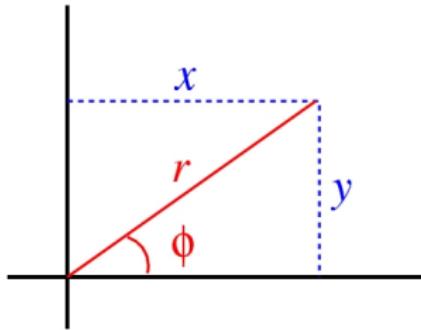
## Model



- Cartesian coord.  $(x, y)$

$$x^2 + y^2 = R^2$$

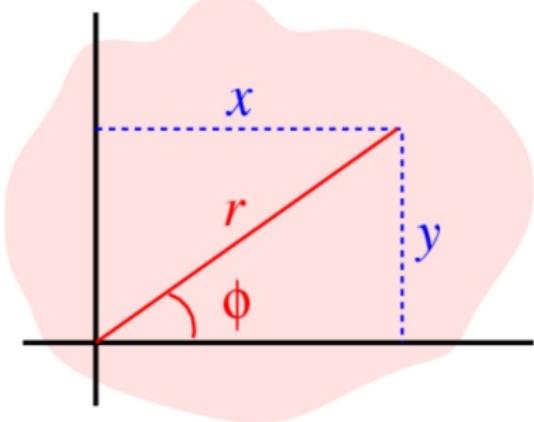
- Polar coord.  $(r, \phi)$





# Polar coordinates

- Transformations



$$r = \sqrt{x^2 + y^2} \quad \phi = \arctan\left(\frac{y}{x}\right)$$

$$y = r \sin \phi \quad x = r \cos \phi$$

- Intervals

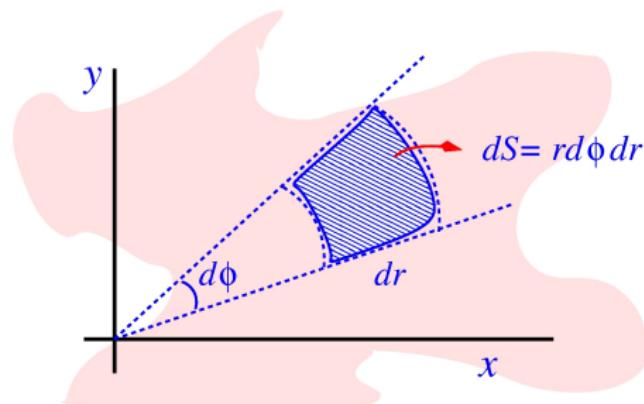
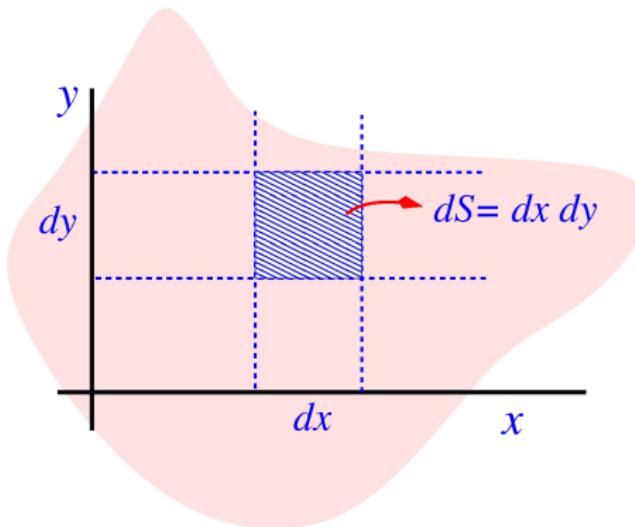
$$r \in [0, +\infty)$$

$$\phi \in [0, 2\pi)$$



## Polar coordinates

- Differential element of surface  $\Rightarrow dS = r dr d\phi$



$$\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} f(x, y) dx dy = \int_0^{+\infty} \int_0^{2\pi} f(r, \phi) r dr d\phi$$

$$\hat{H}\Psi = E\Psi$$

# Eigenfunctions and eigenvalues

- Hamiltonian in polar coordinates

$$\hat{H}(x, y) = -\frac{\hbar^2}{2m} \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right)$$

$$\hat{H}(r, \phi) = -\frac{\hbar^2}{2m} \left( \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \right)$$

$$\hat{H}(\phi) = -\underbrace{\frac{\hbar^2}{2mR^2}}_I \frac{\partial^2}{\partial \phi^2}$$

- Schrödinger's equation

$$-\frac{\hbar^2}{2I} \frac{\partial^2 \psi(\phi)}{\partial \phi^2} = E\psi(\phi)$$

$$\hat{H}\Psi = E\Psi$$

# Eigenfunctions and eigenvalues

- Independent solutions

$$\psi_1(\phi) = e^{k\phi}$$

$$\psi_2(\phi) = e^{-k\phi}$$

$$\frac{\partial^2 \psi_1}{\partial \phi^2} = k^2 \psi_1$$

$$\frac{\partial^2 \psi_2}{\partial \phi^2} = k^2 \psi_2$$

$$k = \sqrt{\frac{2IE}{\hbar^2}} i$$

$$\psi(\phi) = A e^{im\phi} + B e^{-im\phi}$$

$$m = \sqrt{\frac{2IE}{\hbar^2}}$$

$$\text{H}\Psi = E\Psi$$

## Degeneracy

- Single valued function  $\Rightarrow \psi(\phi) = \psi(\phi + 2\pi)$

$$A e^{im\phi} + B e^{-im\phi} = A e^{im(\phi+2\pi)} + B e^{-im(\phi+2\pi)}$$

$$A e^{im\phi} (1 - e^{2\pi m i}) + B e^{-m\phi i} (1 - e^{-2\pi m i}) = 0$$

$$\Rightarrow 1 - e^{-2\pi m i} = e^{-2\pi m i} (e^{2\pi m i} - 1) = -e^{-2\pi m i} (1 - e^{2\pi m i})$$

$$A e^{im\phi} (1 - e^{2\pi m i}) - B e^{-im\phi} e^{-2\pi m i} (1 - e^{2\pi m i}) = 0$$

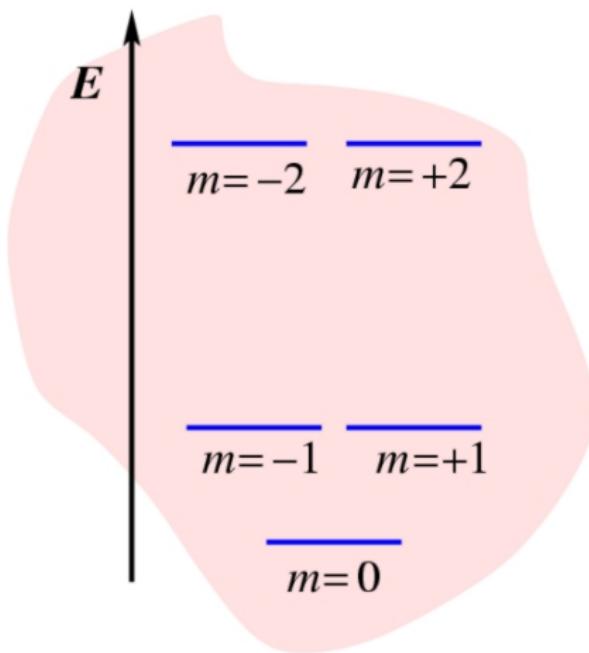
$$\underbrace{(1 - e^{2\pi m i})}_0 (A e^{im\phi} - B e^{-im\phi} e^{-2\pi m i}) = 0$$

$$\cos(2m\pi) + i \sin(2m\pi) = 1 \Rightarrow m = 0, \pm 1, \pm 2, \dots$$

$$\hat{H}\Psi = E\Psi$$

# Degeneracy

- Energy quantization  $\Rightarrow E_m = \frac{\hbar^2 m^2}{2I}$



- Degeneracy

$$E_{-1} = E_1$$

$$E_{-2} = E_2$$

.....



## Angular momentum

- Angular momentum operator  $\Rightarrow$  If particle moves in the  $xy$  plane then  $L = l_z$

$$l_z = xp_y - yp_x \rightarrow \hat{l}_z = \frac{\hbar}{i} \left( x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) = \rightarrow \hat{l}_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$$

- Eigenfunctions and eigenvalues

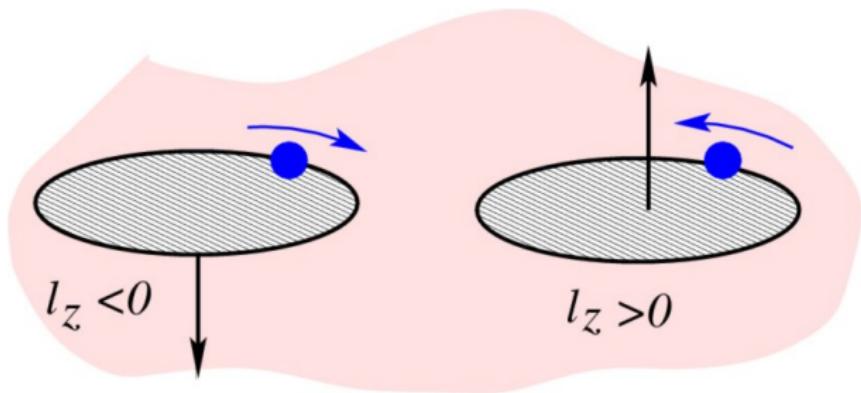
$$\begin{aligned} \hat{l}_z \psi(\phi) &= \frac{\hbar}{i} \frac{\partial}{\partial \phi} \left( A e^{im\phi} + B e^{-im\phi} \right) = m\hbar \left( A e^{im\phi} - B e^{-im\phi} \right) \\ &= l_z \left( A e^{im\phi} + B e^{-im\phi} \right) \end{aligned}$$

$$B = 0 \Rightarrow l_z = m\hbar$$

$$\psi(\phi) = A e^{im\phi} \quad m = 0, \pm 1, \pm 2 \dots$$



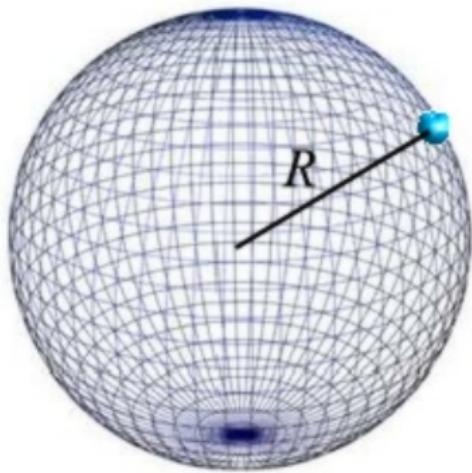
## Angular momentum



- Normalization  $\Rightarrow \int_0^{2\pi} |\psi(\phi)|^2 d\phi = 1 \Rightarrow |A| = \frac{1}{\sqrt{2\pi}}$
- Probability density  $\Rightarrow \psi(\phi) = \frac{1}{\sqrt{2\pi}} e^{im\phi} \Rightarrow |\psi(\phi)|^2 = \frac{1}{2\pi}$



## Model



- Free motion on the surface

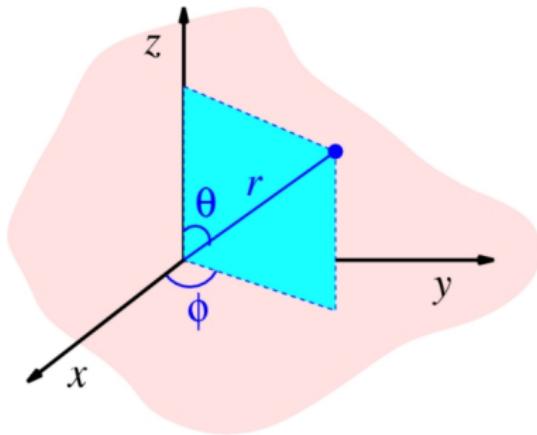
$$V(x, y, z) = 0$$

- Simple atomic model
- Cartesian coordinates

$$x^2 + y^2 + z^2 = R^2$$



## Spherical polar coordinates



- Spherical polar coordinates

$$\begin{aligned}x &= r \sin \theta \cos \phi \\y &= r \sin \theta \sin \phi \\z &= r \cos \theta\end{aligned}$$

- Intervals  $r \in [0, \infty)$ ,  $\phi \in [0, 2\pi)$ ,  $\theta \in [0, \pi)$

- Differential element of volume  $\Rightarrow d\tau = r^2 \sin \theta dr d\theta d\phi$

$$\iiint_{-\infty}^{+\infty} f(x, y, z) dx dy dz = \int_0^{+\infty} \int_0^{\pi} \int_0^{2\pi} f(r, \theta, \phi) r^2 \sin \theta dr d\theta d\phi$$



# Energy and angular momentum

- Angular momentum

$$\hat{L}^2(\theta, \phi) = -\hbar^2 \left( \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} \right)$$

$$\hat{L}^2 Y_l^m(\theta, \phi) = L^2 Y_l^m(\theta, \phi) \rightarrow L^2 = l(l+1)\hbar^2 \rightarrow l = 0, 1, 2, \dots$$

- Energy  $\Rightarrow \hat{H} = \underbrace{\frac{\hat{L}^2}{2mR^2}}_l \Rightarrow E = l(l+1)\frac{\hbar^2}{2I}$

- $z$  component angular momentum

$$\hat{L}_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi} \rightarrow \hat{L}_z Y_l^m(\theta, \phi) = L_z Y_l^m(\theta, \phi)$$

$$L_z = m\hbar \quad m = -l, -l+1, \dots, 0, \dots, l-1, l$$



# Spherical Harmonics

- Spherical Harmonics  $\Rightarrow Y_l^m(\theta, \phi) = \Theta_{l,m}(\theta)\Phi_m(\phi)$

$$\Phi_m(\phi) = \frac{1}{\sqrt{2\pi}} e^{im\phi}$$

$\Theta_{l,m}(\theta) \Rightarrow$  Associated Legendre's functions

| $l$     | $m$                 | $Y_l^m(\theta, \phi)$       |
|---------|---------------------|-----------------------------|
| 0       | 0                   | $1/\sqrt{4\pi}$             |
| 1       | 0                   | $\sqrt{3/4\pi} \cos \theta$ |
| $\pm 1$ | $\pm \sqrt{3/8\pi}$ | $\sin \theta e^{\pm i\phi}$ |



# Probability density



$l=0, (m=0)$



$l=1, (m=0, |m|=1)$



$l=2, (m=0, |m|=1, |m|=2)$

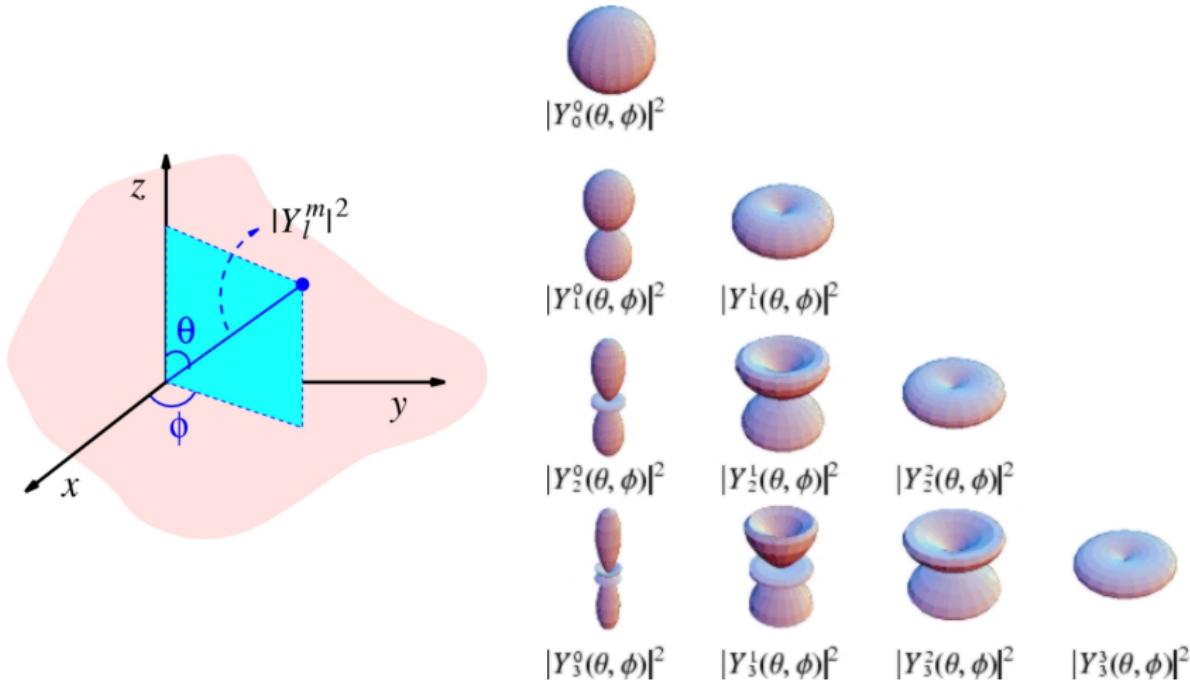


$l=3, (m=0, |m|=1, |m|=2, |m|=3)$





# Polar representations





## Notation

| $l$                                     | Symbol |                    |
|---|--------|--------------------|
| 0                                       | s      | sharp              |
| 1                                       | p      | principal          |
| 2                                       | d      | diffuse            |
| 3                                       | f      | fundamental        |
| 4                                       | g      | alphabetical order |
| :                                       | :      | :                  |
| $m \Rightarrow$ subscript if $l \neq 0$ |        |                    |

### • Examples

$$Y_0^0 \rightarrow s$$

$$Y_1^0 \rightarrow p_0$$

$$Y_1^1 \rightarrow p_1$$

$$Y_1^{-1} \rightarrow p_{-1}$$

$$Y_2^0 \rightarrow d_0$$

$$Y_2^2 \rightarrow d_2$$



$\hat{L}_x, \hat{L}_y?$

- $[\hat{L}_x, \hat{L}_z] \neq 0, [\hat{L}_y, \hat{L}_z] \neq 0$

